

# Low-dimensional magnetism by Dipankar Sen

Ferromagnets & Antiferromagnets



$$H = -J \sum_n \vec{S}_n \cdot \vec{S}_{n+1}$$

$$S_n^2 = S(S+1) \hbar^2$$

$J > 0$

Ferromagnetic system



All  $S_n^z = S \hbar$

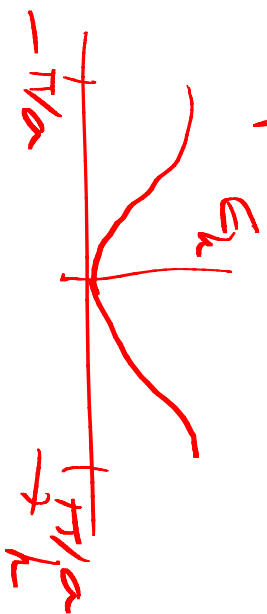
Spin-wave theory:

Excited state with momentum  $k$

$$|k\rangle = \sum_n e^{ikn} \begin{array}{cccccccc} & \uparrow & \uparrow & \uparrow & \dots & \uparrow & \uparrow & \uparrow & \uparrow & \uparrow \\ S & S & S & S & \dots & S & S & S & S & S \\ & s-1 & n & s-1 & & n & s-1 & s-1 & s-1 & s-1 \end{array}$$

$$E_k = 2JS (1 - \cos(ka))$$

Spin-wave (magnon) dispersion



Gapless at  $k=0$

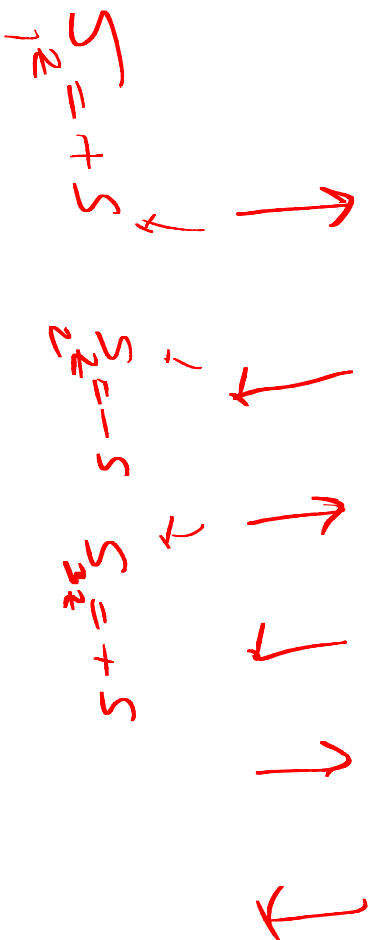
Around  $k=0$ ,

$$E_k \sim k^2$$

Anti ferromagnet

$$H = J \sum_n \vec{S}_n \cdot \vec{S}_{n+1}$$

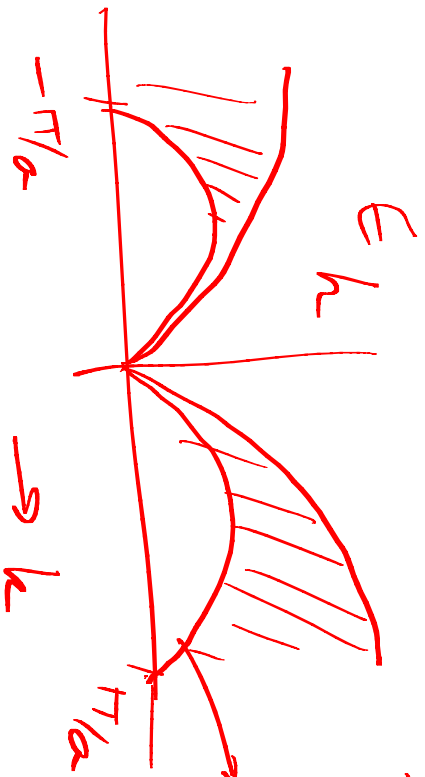
$$J > 0$$



Neel configuration

Bethe ansatz : Ground state has to be spin  
 lowest excited state  
 ( $S=0$  and  $1$ )  
 (singlet)

Spin- $\frac{1}{2}$   
 AFM  
 chain



$$\pi J \sin\left(\frac{ka}{2}\right)$$

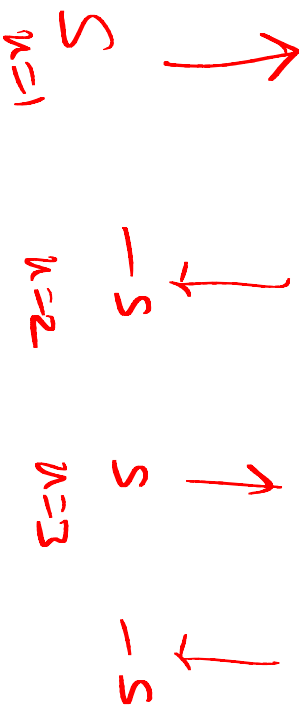
$$\frac{\pi J}{2} \sin(ka)$$

Gaps at  $\pm \frac{\pi}{a}$   
 and  $k=0$   
 Dispersion is linear  
 at  $k=0, \pi$

$$S = 1, 3, \dots ?$$

$S \rightarrow \infty$  Semiclassical limit

Andersen, Phys Rev. 86, 694 (1952)



Holstein-Primakoff  
transf.

$$\begin{aligned}
 n=1: \quad S_{n2} &= S - a_n^\dagger a_n \\
 S_{nt} &\approx \sqrt{2S} a_n, \quad S_{nt} \approx \sqrt{2S} a_n^\dagger
 \end{aligned}$$

$n=2$ :

$$S_{n+} = -S + a_n^T a_n$$

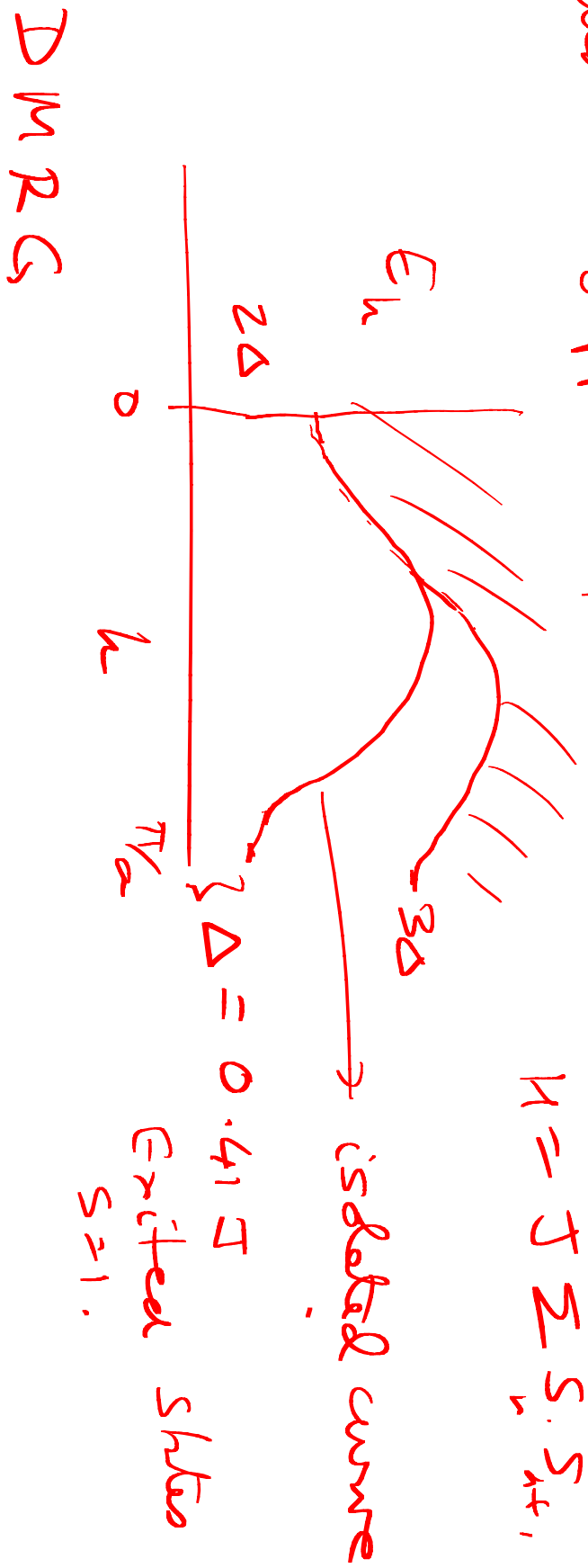
$$S_{n+} = \sqrt{2S} a_n^T, \quad S_{n-} = \sqrt{2S} a_n$$

Write  $H$  in terms of  $n_n, a_n^T$  up to second order.  
 $\sim a^T a, a^T a^T, a a$



Madhava (1981-83) that the spectrum is gapped only if  $S = 1/2, 3/2, 5/2, \dots$  and is gapped if  $S = 1, 2, 3, \dots$

$S = 1/2, 3/2, 5/2, \dots$  → gapped. States  $\psi_{k+1/2}$  and  $\psi_{k+3/2}$  are non-degenerate  
 $S = 1, 2, 3, \dots$  →  $H = J \sum S_i \cdot S_{i+1}$



For general integer  $S$ ,  
 $\Delta \sim e^{-\pi S}$  (Haldane, Affleck...)  
 $\rightarrow 0$  as  $S \rightarrow \infty$ .

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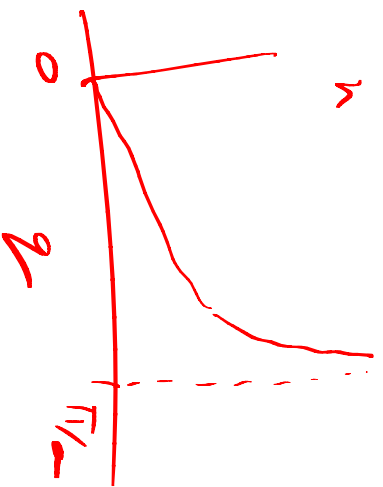
Structure function:  $S(g)$

$$\langle \vec{S}_0 \cdot \vec{S}_n \rangle \underset{\text{ground state}}{\sim} \frac{(-1)^n}{|n|} \sqrt{\ln |n|} \quad S = \frac{1}{2}$$

$\rightarrow e^{i\pi n}$



$$S(q) = \sum_n e^{iqna} \langle \vec{S}_0 \cdot \vec{S}_n \rangle \xrightarrow{\text{quad. str.}} S(q) \rightarrow \infty \text{ as } 2 \rightarrow \pi/a$$

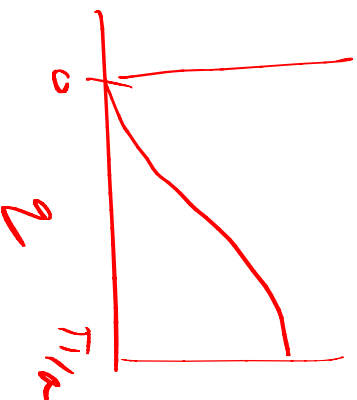


$$\sum_n \vec{S}_n = 0 \text{ in quad. str.}$$

$$S = 1: \quad \langle S_0 \cdot S_n \rangle \sim \frac{(-1)^n}{|n|} \quad e^{-n/\xi} \quad \xi \sim \frac{1}{\Delta}$$

$$\xi = \text{correlation length} \approx 6$$

$S(q)$

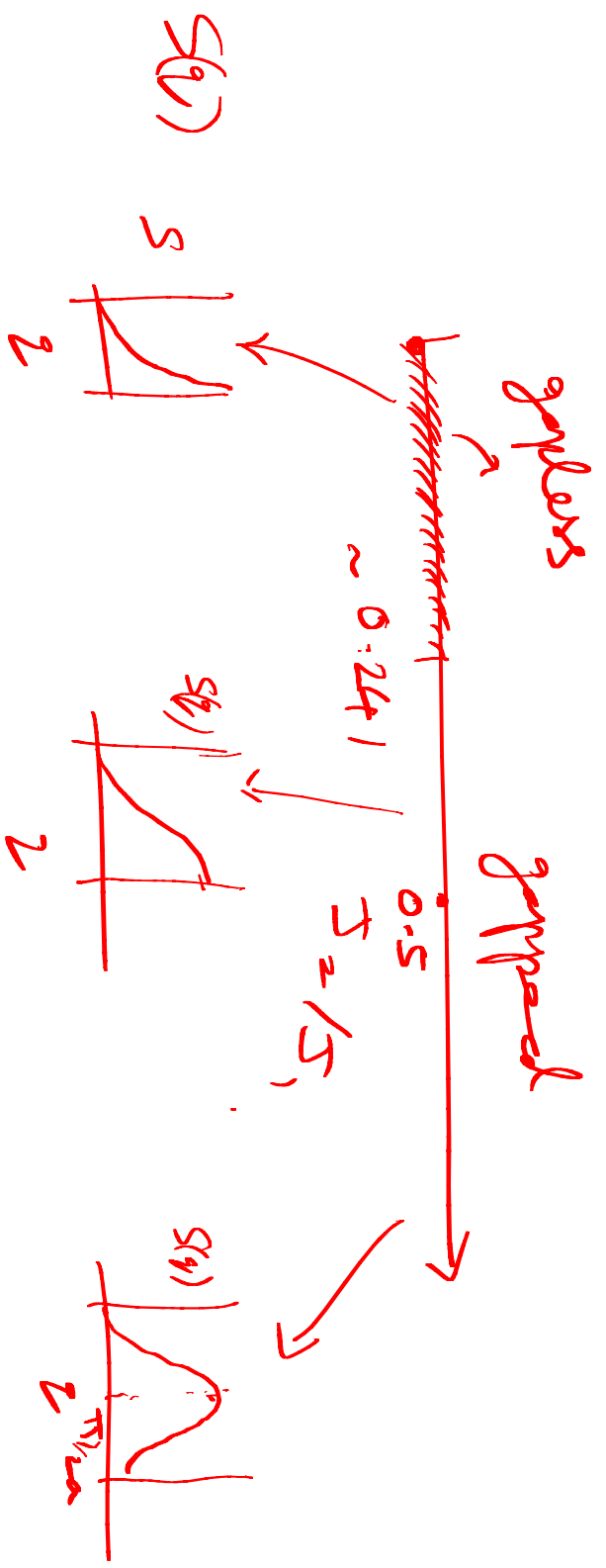
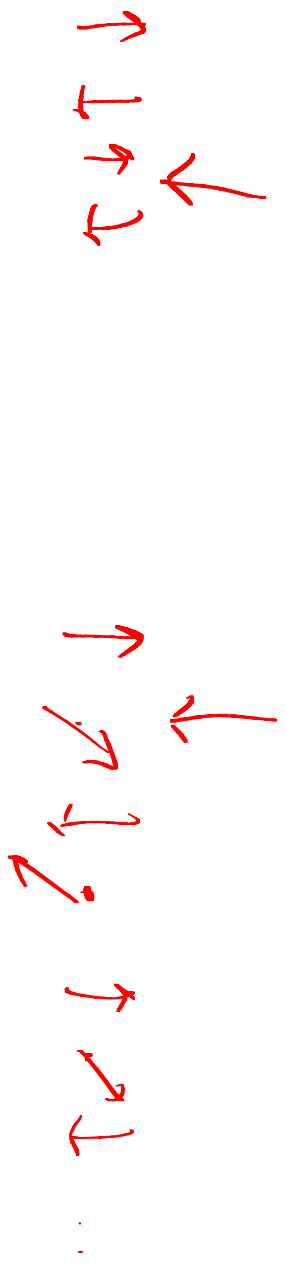


Frustrated model:

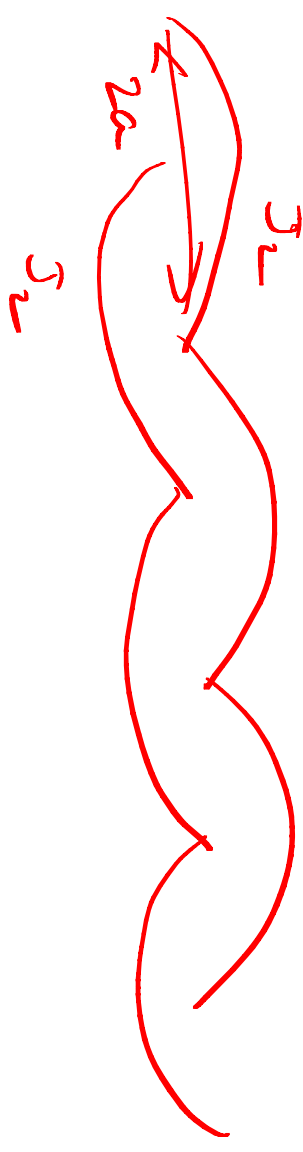
$$H = J_1 \sum_n \vec{S}_n \cdot \vec{S}_{n+1} + J_2 \sum_n \vec{S}_n \cdot \vec{S}_{n+2}$$

Spin  $\frac{1}{2}$  model

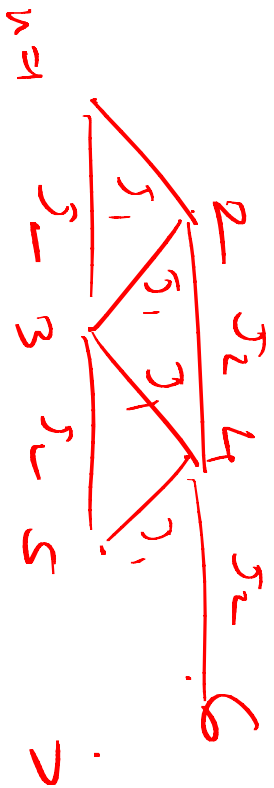
$J_2 \rightarrow \infty$ :  
Two decoupled chains



$$H = J_1 \sum S_x^i \cdot S_x^{i+1} + J_2 \sum S_x^i \cdot S_x^{i+2}$$



$$J_1 \sim 0$$



$S_2/T_1 = 1/2$  ; Maximum - Share in all

$$\begin{aligned}
 H &= S_1 \left[ \sum_{i=1}^n S_i^{-1} \cdot S_{n+i}^{-1} + \frac{1}{2} \sum_{i=1}^n S_i^{-1} \cdot S_{n+i}^{-1} \right] \\
 &= S_1^{-1} \sum_{i=1}^n (S_n^{-1} + S_{n+i}^{-1} + S_{n+i}^{-2})^2 + \text{constant.}
 \end{aligned}$$

Grnd. slts is a product of singlets over nearest neighbors

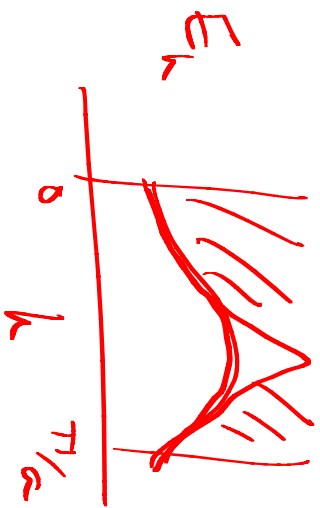


$$\prod_n \frac{1}{\sqrt{L}} \left( \uparrow, \downarrow, \downarrow, \dots - \downarrow, \uparrow, \uparrow, \dots \right)$$

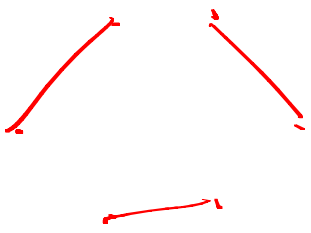
Grnd. slts  
S=0



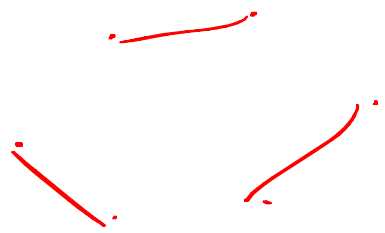
$$\prod_n \frac{1}{\sqrt{L}} \left( \uparrow_0, \downarrow_1, \downarrow_2, \dots - \downarrow_0, \uparrow_1, \uparrow_2, \dots \right)$$



Benzene



$n$



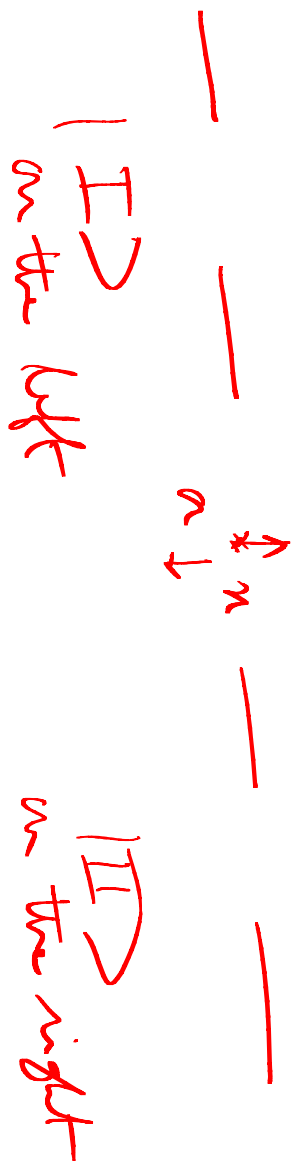
Slonitsky & Sutherland,  
 Phys. Rev. Lett.  
 42, 964 (1981).



$|II\rangle =$

Soliton  
Larmor walk

Lowest  
possible  
excitation



$|II\rangle$   
on the right

Spinon

$$\rightarrow |k\rangle = \sum_n e^{ikn} |n\rangle$$



Dimerisation



(Cu<sup>2+</sup> = Sp<sup>n-1/2</sup>)

$$s \approx 0.06$$

$$\vec{s}_1 \quad \vec{s}_2 \quad \vec{s}_3 \quad s_4 \quad \dots$$

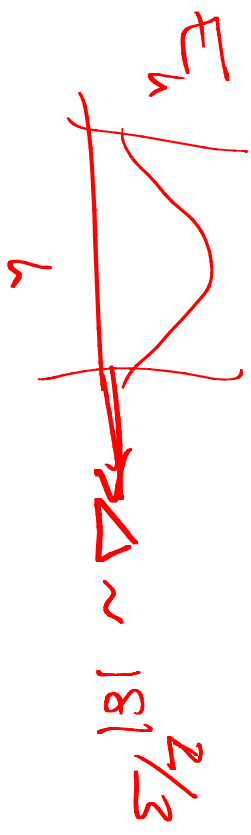
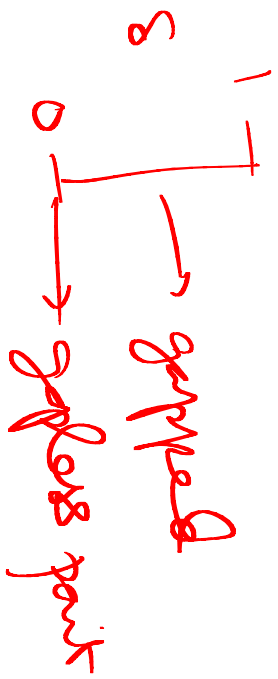
$$J_1(1+s) \quad J_1(1-s) \quad J_1(1+s) \quad J_1(1-s)$$

→ Periodic  
in S hybridity

$$H = J_1(1+s) \vec{s}_1 \cdot \vec{s}_2 + J_1(1-s) \vec{s}_2 \cdot \vec{s}_3 + \dots$$

$$s \rightarrow 1: \quad \frac{1}{2s_1} \quad 0 \quad \frac{1}{2s_1} \quad 0 \quad \dots$$

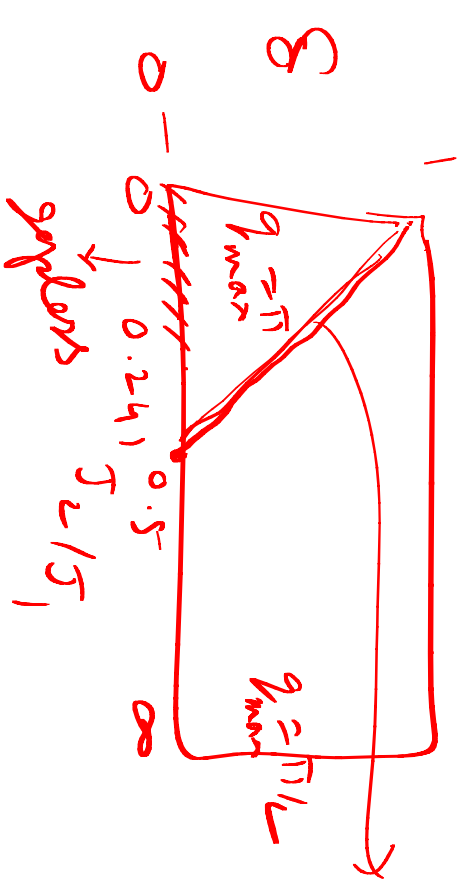
Grid. site is a singlet and non-degenerate



$$\chi \sim e^{-\Delta/k_B T} \text{ as } T \rightarrow 0$$

$\sim$  constant as  $T \rightarrow 0$   
for gapless systems -

$$H = J_1 \sum_n [1 + (-1)^n \delta] \vec{S}_n \cdot \vec{S}_{n+1} + J_2 \sum_n \vec{S}_n \cdot \vec{S}_{n+2}$$



$\frac{2J_2}{J_1} + \delta = 1$   
 $\rightarrow$  ground state exactly  
 Schrodinger  
 $S(q) \rightarrow q_{max}$

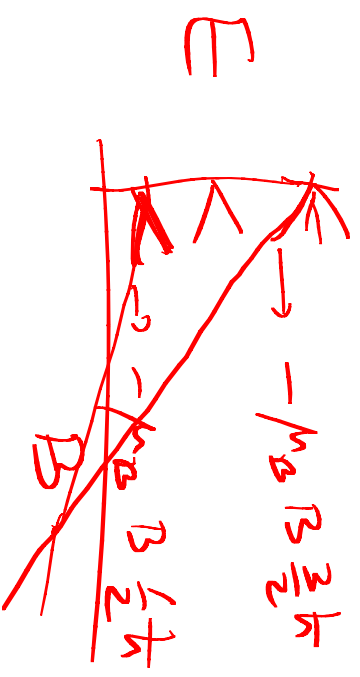
# Magnetization plateaus

2:  $H = J_1 (S_1^z \cdot S_2^z + S_2^z \cdot S_3^z)$

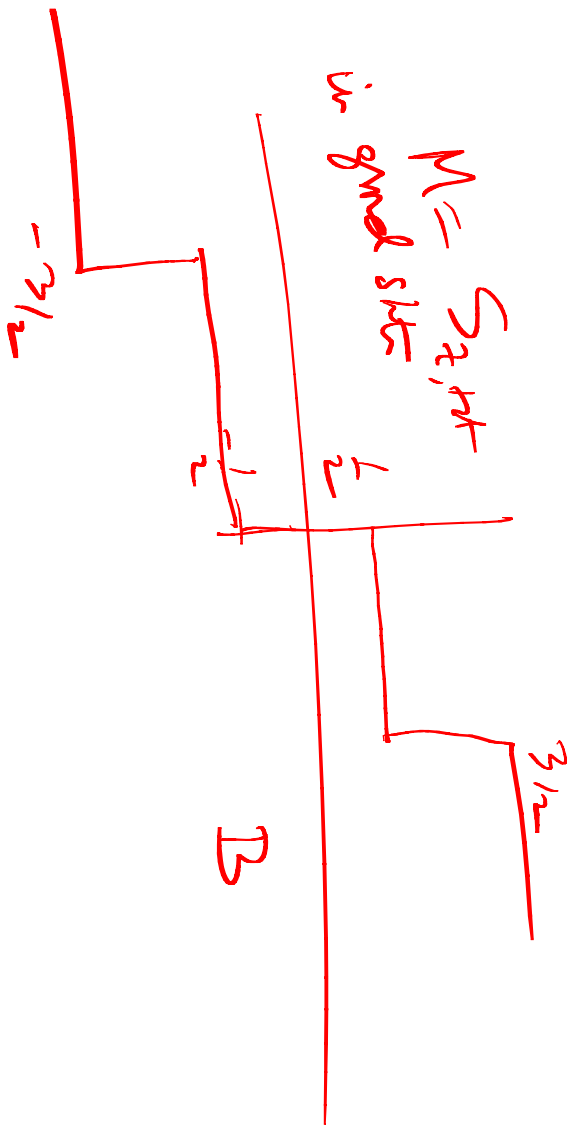
1 + 3  $J_1 > 0$

- $S = 3/2$
- $S = 1/2$
- $S = 1/2$

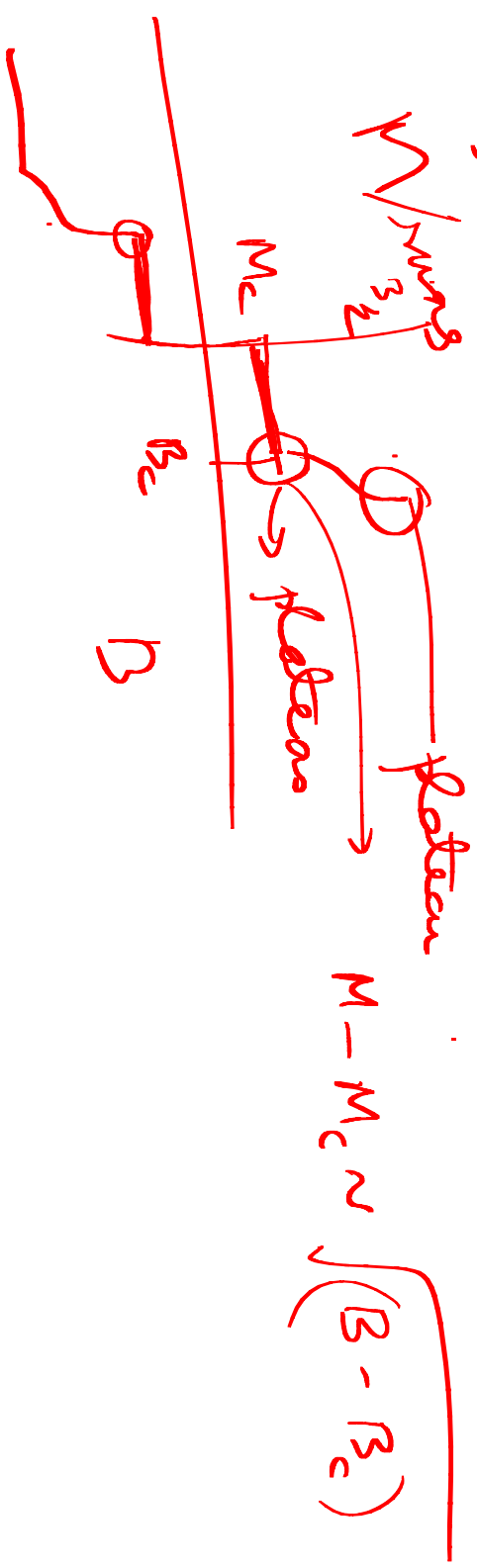
$H = J_1 (S_1^z \cdot S_2^z + S_2^z \cdot S_3^z) - \mu_B B (S_{1z} + S_{2z} + S_{3z})$



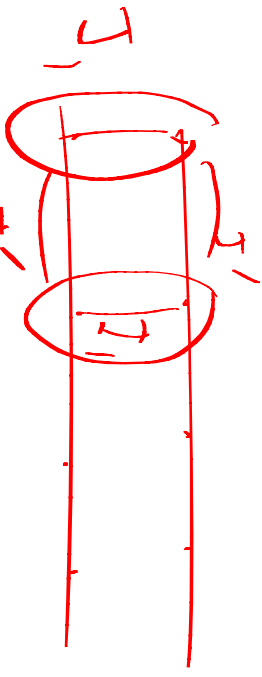
More save B with  $S = 3/2$ ,  $S_z = 3/2$   
because the good site



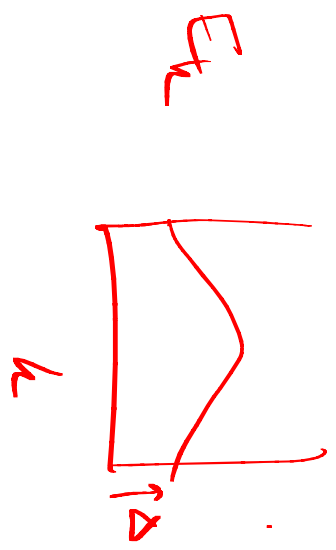
Gebora, Hamecker and Pijzel, Phys. Rev. Lett. 79, 5126 (1977).  
 3-log ladder



$$M - M_c \sim \sqrt{(B - B_c)}$$



$S=0$  in layer spin  $\rightarrow$  Hallane gap



ground stts  $\neq S=0$   
 excited stts  $\Rightarrow S=1$

