

## Introduction to Magnetism Part II

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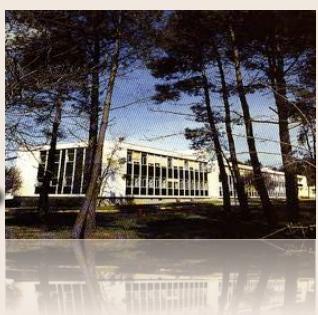
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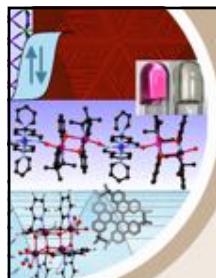


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*Introduction to Magnetism... by R. Clérac*

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## Introduction to Magnetism

- 1) Introduction and definitions**
- 2) The Diamagnetism**
- 3) The Paramagnetism**
- 4) Magnetic phase transitions and magnetic orders**
- 5) Molecular magnetism**
- 6) Superparamagnetism and Single-Molecule Magnets**
- 7) Single-Chain Magnets**

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**6) Superparamagnetism and SMMs**

**Single domain magnetic nanoparticles (with uniaxial anisotropy): classical case of uniform rotation**

$H \neq 0$

$M \neq 0$

$T = 25^\circ\text{C}$

$E (\text{J})$

$\theta (^\circ)$

$\Delta \propto KN$

super para.      single domain

Magnetic Material

Particle Diameter (nm)

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**6) Superparamagnetism and SMMs**

**Single domain magnetic nanoparticles (with uniaxial anisotropy): classical case of uniform rotation**

$M \neq 0$

$M / M_s = \exp\left(-\frac{t}{\tau}\right)$

$t / \tau$

Temperature effect:  $\tau(T) = \tau_0 \exp\left(\frac{KN}{k_B T}\right)$  activated relaxation

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**6) Superparamagnetism and SMMs**

Single domain magnetic nanoparticles (with uniaxial anisotropy): classical case of uniform rotation

$M = 0$

Temperature effect:  $\tau(T) = \tau_0 \exp\left(\frac{KN}{k_B T}\right)$  activated relaxation

bistability, hysteresis when  $\tau_{\text{exp}} < \tau(T)$

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**6) Superparamagnetism and SMMs**

Information storage... Hard Disk Drives (HDD)...

Conventional Multigrain Media

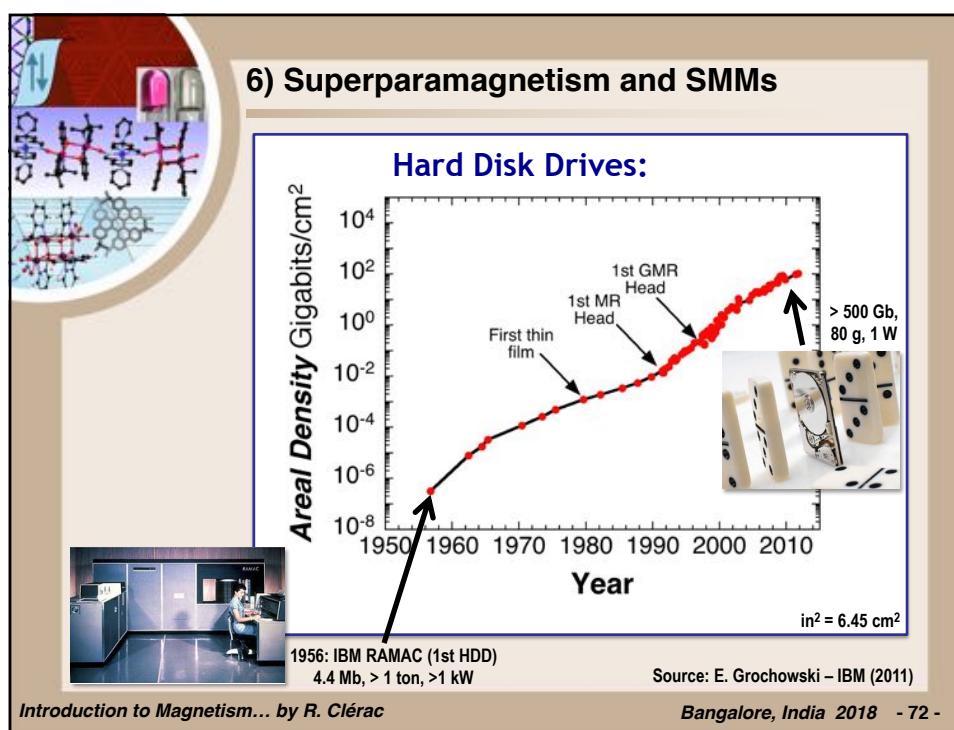
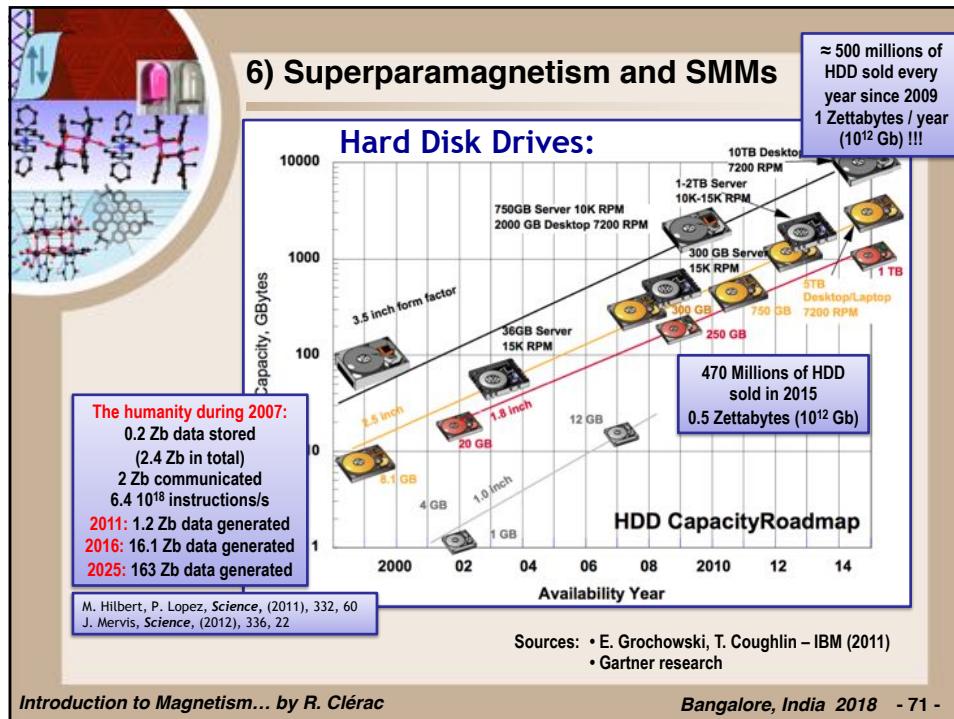
200 nm

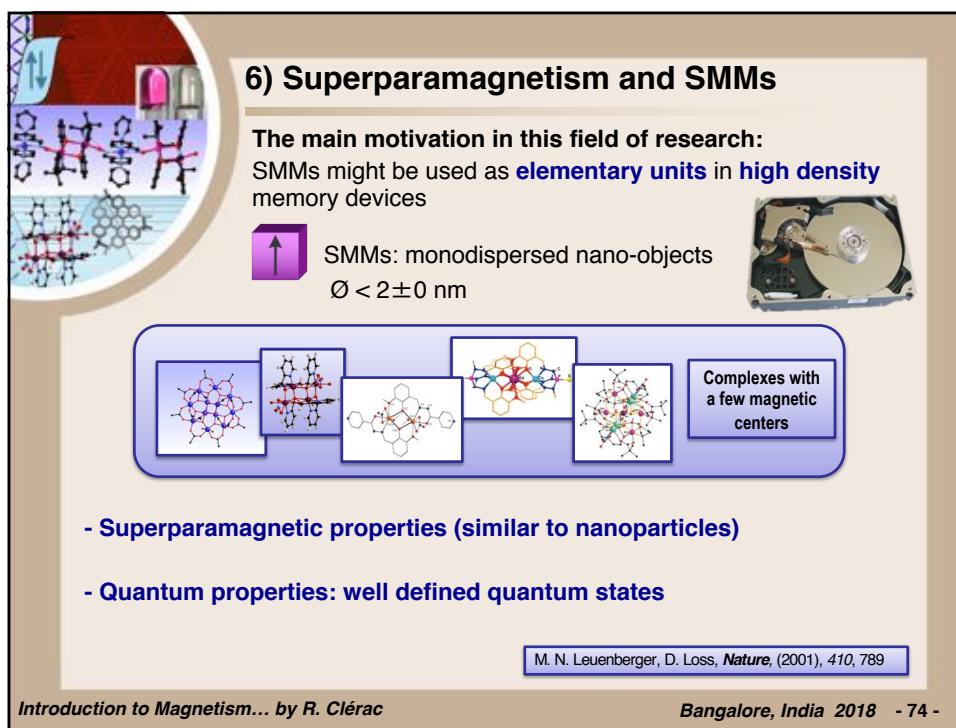
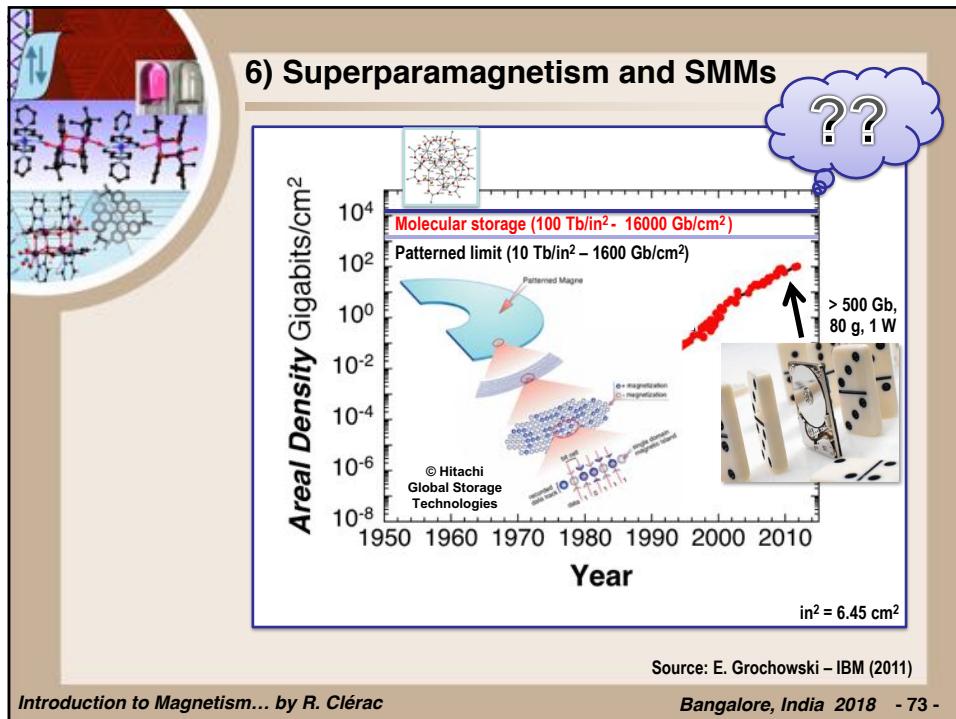
Particle/grain sizes:  $\varnothing = 5-8 \text{ nm}$

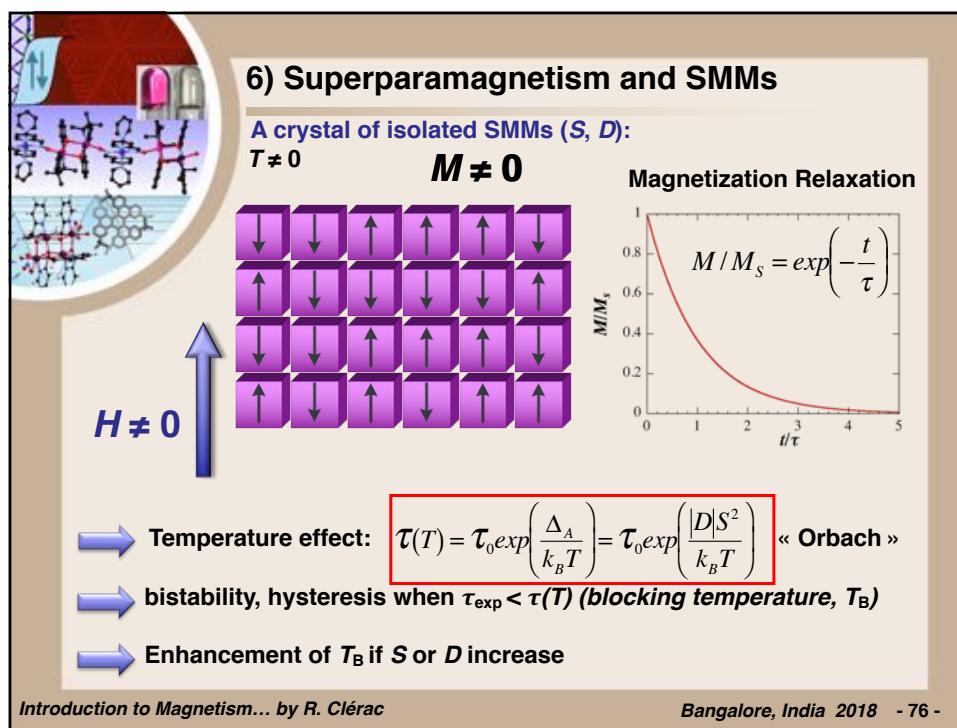
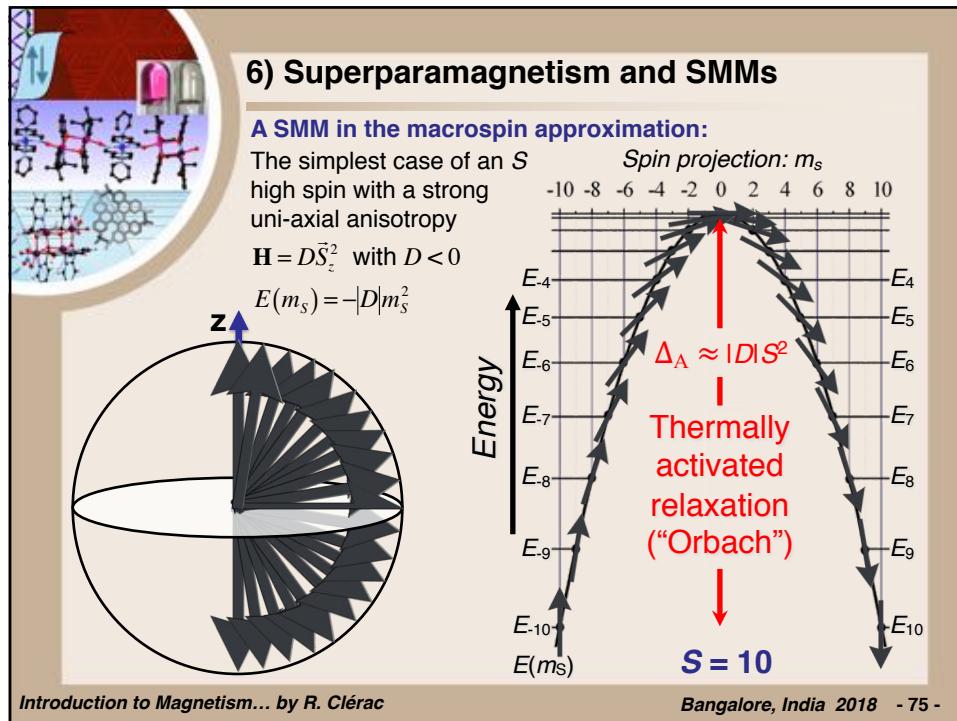
10 nm

IBM Hard Drive (35 Gb/in<sup>2</sup>; CoPtCrB)  
10 years of storage

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**6) Superparamagnetism and SMMs**

Determination of the relaxation time: an example...

**(NEt<sub>4</sub>)<sub>2</sub>[Mn<sup>III</sup><sub>2</sub>(salmen)<sub>2</sub>(MeOH)<sub>2</sub>Fe<sup>III</sup>(CN)<sub>6</sub>]**

salmen<sup>2</sup>: *rac*-*N,N*-(1-methylethylene)bis(salicylideneiminate)

$S = 9/2$   
 $D/k_B = -1.25 \text{ K}$

**Direct measurements in time (dc)**

M. Ferbinteanu, H. Miyasaka, W. Wernsdorfer, K. Nakata, K. Sugiura, M. Yamashita, C. Coulon and R. Clérac, *J. Am. Chem. Soc.* **2005**, *127*, 3090

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**6) Superparamagnetism and SMMs**

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Relaxation time ( $\tau$ ) deduced from the direct measurements in time

$\tau / \text{s}$

$T^{-1} / \text{K}^{-1}$

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**6) Superparamagnetism and SMMs**

Determination of the relaxation time: an example...  
 $(\text{NEt}_4)[\text{Mn}^{\text{III}}_2(\text{salmen})_2(\text{MeOH})_2\text{Fe}^{\text{III}}(\text{CN})_6]$

R. Boca, *Theoretical Foundations of Molecular Magnetism*, Elsevier Ed., p.107.

$S = 9/2$   
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Measurements in frequency (ac)

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Measurements in frequency (ac)

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 $(\text{NEt}_4)[\text{Mn}^{\text{III}}_2(\text{salmen})_2(\text{MeOH})_2\text{Fe}^{\text{III}}(\text{CN})_6]$

Relaxation time ( $\tau$ ) deduced from the dc and ac measurements

$\tau_{\text{QTM}} = 455 \text{ s}$

$S = 9/2$   
 $D/k_B = -1.25 \text{ K}$

→  $\tau = \tau_{\text{QTM}}$  because  $\tau_{\text{Orbach}} > \tau_{\text{QTM}}$  below  $T_{\text{QTM}}$

→ A second mode of relaxation: Quantum Tunneling

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**6) Superparamagnetism and SMMs**

A second mode of relaxation: case for  $k_B T \ll |E_{10} - E_9|$

To have quantum tunneling, the symmetry of the anisotropy must be lower than uni-axial:

$$\mathbf{H} = D\vec{S}_z^2 + E(\vec{S}_x^2 - \vec{S}_y^2)$$

Spin projection:  $m_s$

$\pm m_s$  states becomes quantically mixed

Quantum tunneling

$\tau_{\text{QTM}}$

$S = 10$

$E(m_s)$

$E_{-10}$

$E_0$

$E_1$

$E_2$

$E_3$

$E_4$

$E_5$

$E_6$

$E_7$

$E_8$

$E_9$

$E_{-9}$

$E_{-8}$

$E_{-7}$

$E_{-6}$

$E_{-5}$

$E_{-4}$

$E_{-3}$

$E_{-2}$

$E_{-1}$

$E_0$

$E_1$

$E_2$

$E_3$

$E_4$

$E_5$

$E_6$

$E_7$

$E_8$

$E_9$

$E_{-10}$

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**6) Superparamagnetism and SMMs**

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$T^{-1} / \text{K}^{-1}$

$S = 9/2$

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**6) Superparamagnetism and SMMs**

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Relaxation time ( $\tau$ ) deduced from the dc and ac measurements

$\tau / \text{s}$

$T^{-1} / \text{K}^{-1}$

$\Delta_{\text{eff}}/k_B = 14 \text{ K}$

$S = 9/2$

$D/k_B = -1.25 \text{ K}$

thus:  
 $|ΔS^2|/k_B = 25 \text{ K}$

→  $Δ_{\text{eff}} < Δ_A$

→ A intermediate mode of relaxation that is called “Thermally assisted quantum tunneling of the magnetization”

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**6) Superparamagnetism and SMMs**

**An intermediate mode of relaxation:  $T > T_{QTM}$**

To have quantum tunneling, the symmetry of the anisotropy must be lower than uni-axial:

$$\mathbf{H} = D\vec{S}_z^2 + E(\vec{S}_x^2 - \vec{S}_y^2)$$

**Spin projection:  $m_s$**

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**6) Superparamagnetism and SMMs**

**Determination of the relaxation time: an example...**

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**Relaxation time ( $\tau$ ) deduced from the dc and ac measurements**

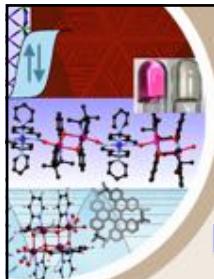
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M. Ferbinteanu, H. Miyasaka, W. Wernsdorfer, K. Nakata, K. Sugiura, M. Yamashita, C. Coulon and R. Clérac, *J. Am. Chem. Soc.* 2005, 127, 3090

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## 6) Superparamagnetism and SMMs

**Beyond QTM and “Orbach” relaxation processes:**

$$\tau^{-1} = \tau_{QTM}^{-1} + \tau_{Orbach}^{-1} + \tau_{direct}^{-1} + \tau_{Raman}^{-1} + \dots$$

R. Orbach Proc. Royal Soc. of London 1961, 264, 458.

Spin-lattice relaxation in rare-earth salts

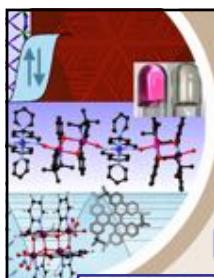
By R. ORBACH†‡  
Clarendon Laboratory, University of Oxford

(Communicated by B. Bleaney, F.R.S.—Received 16 March 1961—  
Revised 26 May 1961)

**1. INTRODUCTION**

The first theoretical treatment of spin-lattice relaxation in solids was given by Waller (1932). He considered in detail modulation of the internal dipolar fields by lattice vibrations, but found very long relaxation times at low temperatures which did not agree with subsequent experiments of Gorter (1936). It remained for Heitler

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R. Orbach Proc. Royal Soc. of London 1961, 264, 458.

For two-phonon processes, the additional distinction of whether the Debye energy ( $K\theta_D$ ) is less than or greater than the crystalline field splitting  $\Delta$  between the ground state and the first excited state must be made. Non-Kramers salts to which the former condition apply ( $K\theta_D < \Delta$ ) are shown to possess two-phonon relaxation processes of the usual Raman type. The relaxation time is proportional to  $T^{-1}$  and is independent of magnetic field. When  $K\theta_D > \Delta$ , there is present in addition a term arising from a resonance process, analogous to the resonance radiation effect in gases. Phonons of energy  $\sim \Delta$  are absorbed and emitted by the spin system preferentially because of a phonon resonance with the crystalline field splitting of the spin states. As normally  $KT$  is much less than  $\Delta$ , this leads to a relaxation time proportional to  $\exp(\Delta/KT)$ . This process will dominate the Raman process except at very high and low temperatures. It is shown to be significant right down to the liquid-helium range by comparison with the relaxation rate due to direct processes.

Kramers salts, when  $K\theta_D < \Delta$ , owing to a cancellation in the rate equation, exhibit a Raman relaxation time proportional to  $T^{-1}$  and independent of field. This ‘Van Vleck cancellation’ is shown to be a consequence of time reversal symmetry. When  $K\theta_D > \Delta$ , the resonance process is also present, the relaxation time again being proportional to  $\exp(\Delta/KT)$ . The resonance process is now shown to be dominant down to 1 or 2°K for many rare-earth salts.

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**6) Superparamagnetism and SMMs**

Beyond QTM and “Orbach” relaxation processes:

$$\tau^{-1} = \tau_{QTM}^{-1} + \tau_{Orbach}^{-1} + \tau_{direct}^{-1} + \tau_{Raman}^{-1} + \dots$$

LA JOURNAL DE PHYSIQUE  
VOLUME 27, NOVEMBRE-DÉCEMBRE 1966, PAGE 782

**A clarification in 1966!!**      **MISE AU POINT**

LA RELAXATION DES SPINS ÉLECTRONIQUES AVEC LE RÉSEAU  
(Théorie élémentaire et méthodes de mesure du temps  $T_1$ )

PAR JACQUES PESCIA,  
Laboratoire de Physique du Solide Faculté des Sciences de Toulouse (1).

**Résumé.** — Dans une première partie, on définit le phénomène de relaxation, puis on précise les notions de température de spins et de réseau. On décrit ensuite les principaux processus de relaxation et l'on explicite le calcul des Temps  $T_1$ .  
Dans une seconde partie, après avoir indiqué les ordres de grandeur des Temps de relaxation à mesurer, et l'origine des procédés de mesure, on passe en revue les techniques actuelles. Pour chaque méthode de mesure, on rappelle le principe sur lequel est fondée cette mesure, puis on décrit l'appareillage utilisé pour sa réalisation et l'on discute les performances obtenues.

**Abstract.** — In the first part, the relaxation phenomenon is described and the question of the existence of spin temperature is discussed. Calculation of spin-lattice relaxation time is explained for the different processes.  
In the second part, we list the various measurement techniques. For each method, the basic principle is brought out, and a description of the apparatus is presented.

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**6) Superparamagnetism and SMMs**

Beyond QTM and “Orbach” relaxation processes:

$$\tau^{-1} = \tau_{QTM}^{-1} + \tau_{Orbach}^{-1} + \tau_{direct}^{-1} + \tau_{Raman}^{-1} + \dots$$

Always a competition between these relaxation processes  
Don't forget that only one relaxation time is measured !!!

<b>Relaxation rate (<math>\tau^{-1}</math>)</b> $\tau_{QTM}^{-1} = \frac{B_1}{1 + B_2 H^2}$ $\tau_{Orbach}^{-1} = \tau_0^{-1} \exp\left(-\frac{\Delta_A}{k_B T}\right)$	<b>Relaxation time (<math>\tau</math>)</b> Temperature independent Increases in $H^2$ Thermally activated dc-field independent (low fields)
<b>Spins</b>	<b>Phonons</b>

J. Pescia J. Phys. 1966, 27, 782.

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**6) Superparamagnetism and SMMs**

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Relaxation rate ( $\tau^{-1}$ )	Relaxation time ( $\tau$ )
$\tau_{Direct}^{-1} = AH^mT$	Linear in temperature Decreases in $H^m$ ( $m = 2$ non-Kramers or $m = 4$ for Kramers)
Spins      Phonons	

J. Pescia J. Phys. 1966, 27, 782.

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**6) Superparamagnetism and SMMs**

Beyond QTM and “Orbach” relaxation processes:

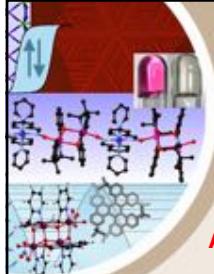
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Relaxation rate ( $\tau^{-1}$ )	Relaxation time ( $\tau$ )
$\tau_{Raman}^{-1} = C \frac{1+C_1 H^2}{1+C_2 H^2} T^n$	A power law in $T$ ( $n = 2-9$ ) Increases or decreases in $H^2$
Spins      Phonons	

J. H. van Vleck, Phys. Rev. 1940, 57, 426.

J. Pescia J. Phys. 1966, 27, 782.

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## 6) Superparamagnetism and SMMs

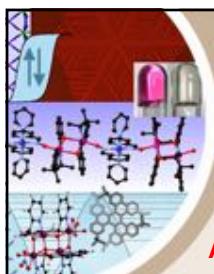
**Beyond QTM and “Orbach” relaxation processes:**

$$\tau^{-1} = \tau_{QTM}^{-1} + \tau_{Orbach}^{-1} + \tau_{direct}^{-1} + \tau_{Raman}^{-1} + \dots$$

A total of 10 parameters if all the relaxations are effective!!!!

Relaxation rate ( $\tau^{-1}$ )	Relaxation time ( $\tau$ )
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$\tau_{Orbach}^{-1} = \tau_0^{-1} \exp\left(-\frac{\Delta_A}{k_B T}\right)$	Thermally activated dc-field independent (low fields)
$\tau_{Direct}^{-1} = A H^m T$	Linear in temperature Decreases in $H^m$ ( $m = 2$ non-Kramers or $m = 4$ for Kramers)
$\tau_{Raman}^{-1} = C \frac{1 + C_1 H^2}{1 + C_2 H^2} T^n$	A power law in $T$ ( $n = 2-9$ ) Increases or decreases in $H^2$

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## 6) Superparamagnetism and SMMs

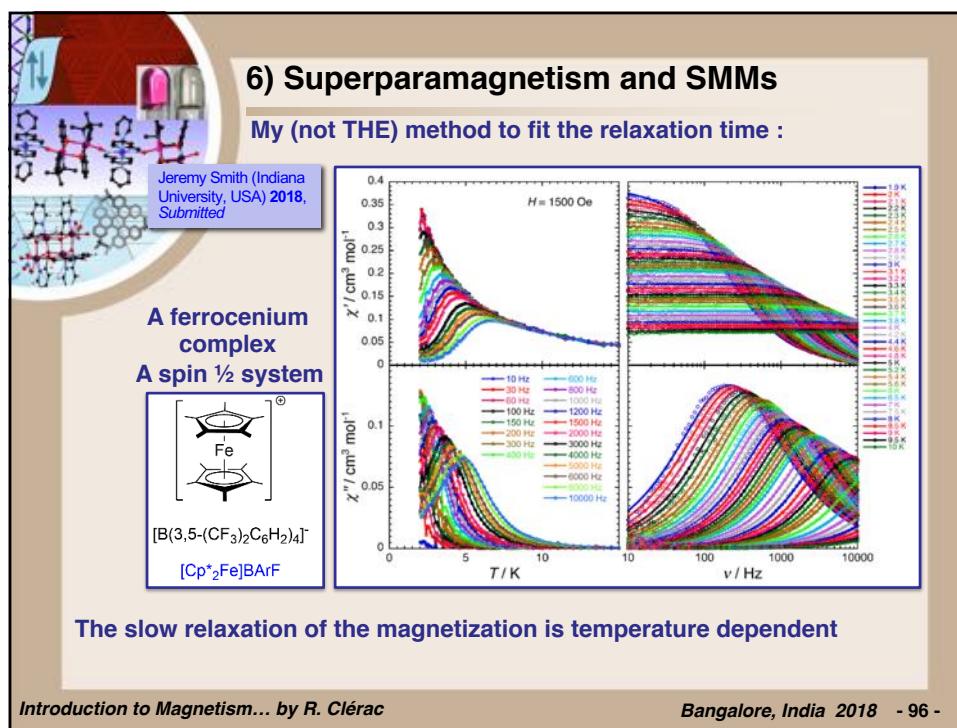
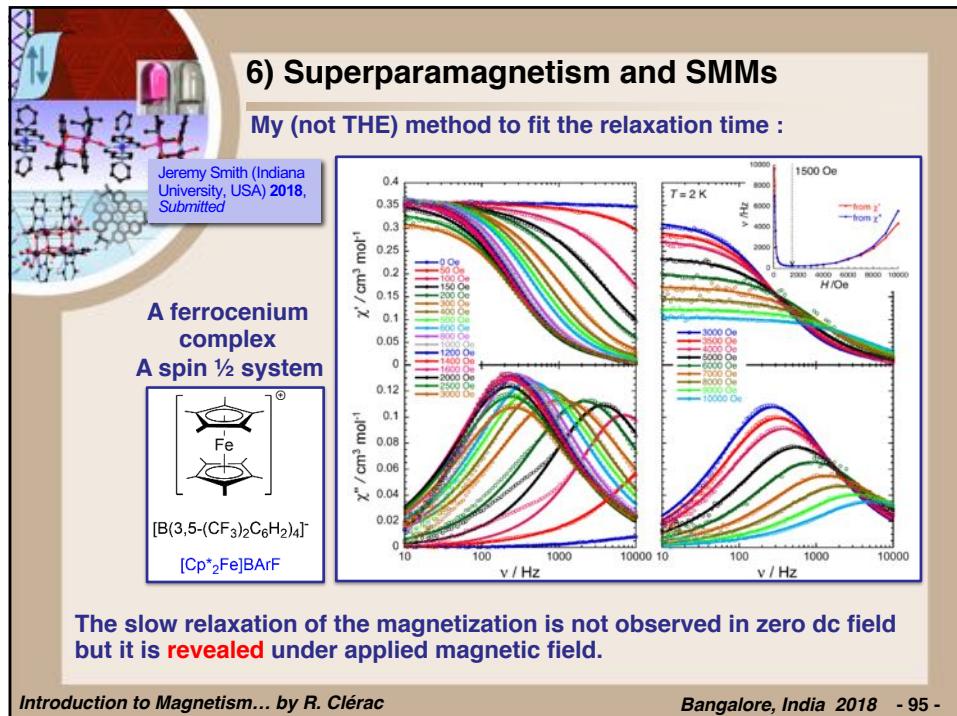
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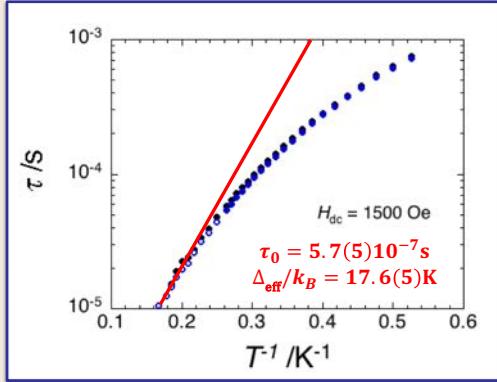
**6) Superparamagnetism and SMMs**

My (not THE) method to fit the relaxation time :

A ferrocenium complex  
A spin ½ system

$\left[ \begin{array}{c} \text{Cp}_2\text{Fe} \\ \text{Fe} \\ \text{Cp}_2 \end{array} \right]^\ominus$

$[\text{B}(3,5-(\text{CF}_3)_2\text{C}_6\text{H}_2)_4]^-$   
 $[\text{Cp}_2^*\text{Fe}]^-\text{BArF}$



- Spin ½ SMM !! Question: “Orbach” and QTM mechanisms?
- To answer: both field and temperature dependences of the experimental relaxation time are **absolutely necessary** !!

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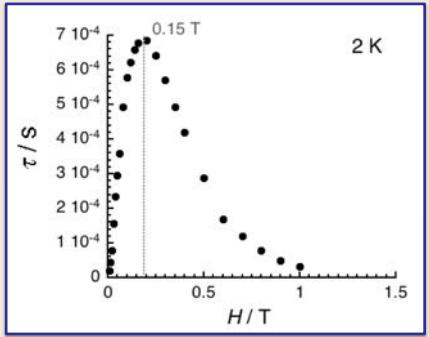
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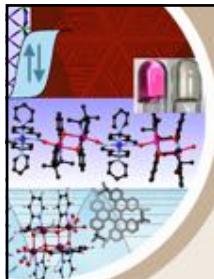
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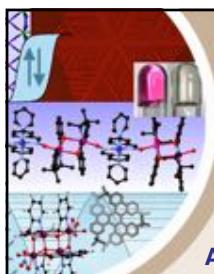
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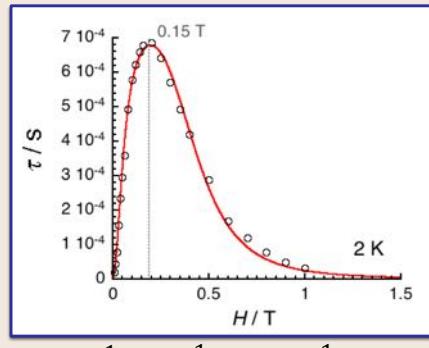
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$\left[ \text{B}(3,5-(\text{CF}_3)_2\text{C}_6\text{H}_2)_4 \right]^\ominus$   
 $[\text{Cp}^*_2\text{Fe}]^+$



$\tau / \text{s}$  vs  $H / \text{T}$ . The plot shows a peak at  $H/T = 0.15 \text{ T}$  and  $T = 2 \text{ K}$ .

$$\tau^{-1} = \tau_{direct}^{-1} + \tau_{Raman}^{-1}$$

$$\tau^{-1} = A H^4 T + C \frac{1 + C_1 H^2}{1 + C_2 H^2} T^n$$

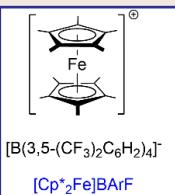
with  $A = 2.0 \cdot 10^4 \text{ K}^{-1} \text{T}^4 \text{s}^{-1}$ ,  $C_1 = 280 \text{ T}^{-2}$ ,  $C_2 = 8.1 \cdot 10^5 \text{ T}^{-2}$  and  $2^n C = 3.8 \cdot 10^6 \text{ s}^{-1}$

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**6) Superparamagnetism and SMMs**

My (not THE) method to fit the relaxation time :

A ferrocenium complex  
A spin 1/2 system



$\tau^{-1} = AH^4T + C \frac{1 + C_1 H^2}{1 + C_2 H^2} T^n$

With only  $n = 3.8$  (optical and acoustic phonons?  $n = 9$  for Kramer systems)  
Fixing  $A = 2.0 \cdot 10^4 \text{ K}^{-1}\text{T}^4\text{s}^{-1}$ ,  $C_1 = 280 \text{ T}^2$ ,  $C_2 = 8.1 \cdot 10^5 \text{ T}^2$  and  $2^n C = 3.8 \cdot 10^6 \text{ s}^{-1}$

K. N. Shriyasta, Phys. Stat. Sol. (b) 1983, 117, 437; A. Singh, K. N. Shriyasta, Phys. Stat. Sol. (b) 1979, 95, 273.

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**6) Superparamagnetism and SMMs**

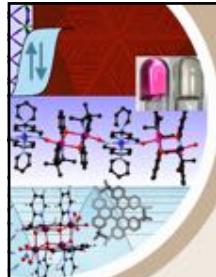
Beyond QTM and Orbach relaxation processes:

$$\tau^{-1} = \tau_{Orbach}^{-1} + \tau_{QTM}^{-1} + \tau_{direct}^{-1} + \tau_{Raman}^{-1}$$

Take home messages:

- All the paramagnetic molecules exhibit slow dynamics of the magnetization!! Known at least since 1932!
- In that sense all the paramagnetic molecules are potentially SMMs!!!
- So what is a SMMs?
- RC: A paramagnetic molecule for which your experiment is able to observe the slow dynamics of the magnetization ( $M$  vs time, ac susceptibility, EPR, NMR, Mossbauer, ...)

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## 6) Superparamagnetism and SMMs

The ingredients in order to obtain a SMM:

- just need a paramagnetic molecule!
- the right experimental setup!!

To promote the “Orbach” process:

- high spin/ magnetic moment
- magnetic anisotropy

Transition metals organized by ligands:

- 3d: Mn(III), Fe(III), Ni(II), Co(II), V(II)
- Lanthanides: Tb(III), Dy(III), Ho(III)
- Mixed metals 3d/3d or 4f/4f or 3d/4f
- Mixed spins: 3d/radicals

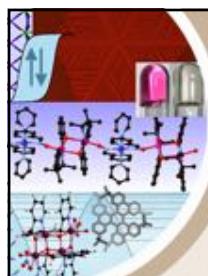
Synthesis methods:

- serendipitous! Most of the time...
- by design (only a few...)

→ There are thousands of SMM examples in the literature

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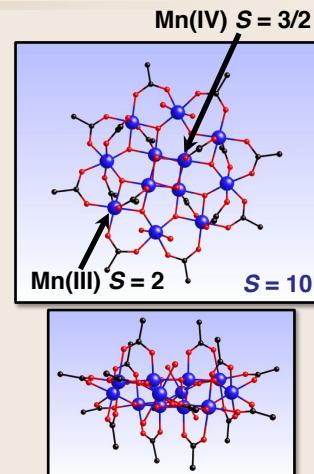
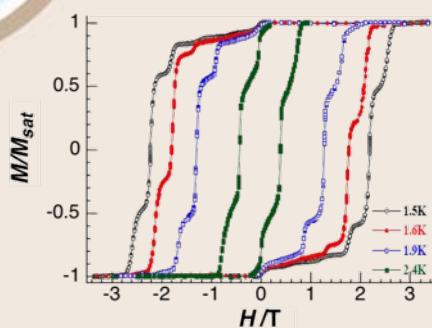
## 6) Superparamagnetism and SMMs

The most famous example!



$$D/k_B = -0.66 \text{ K}$$

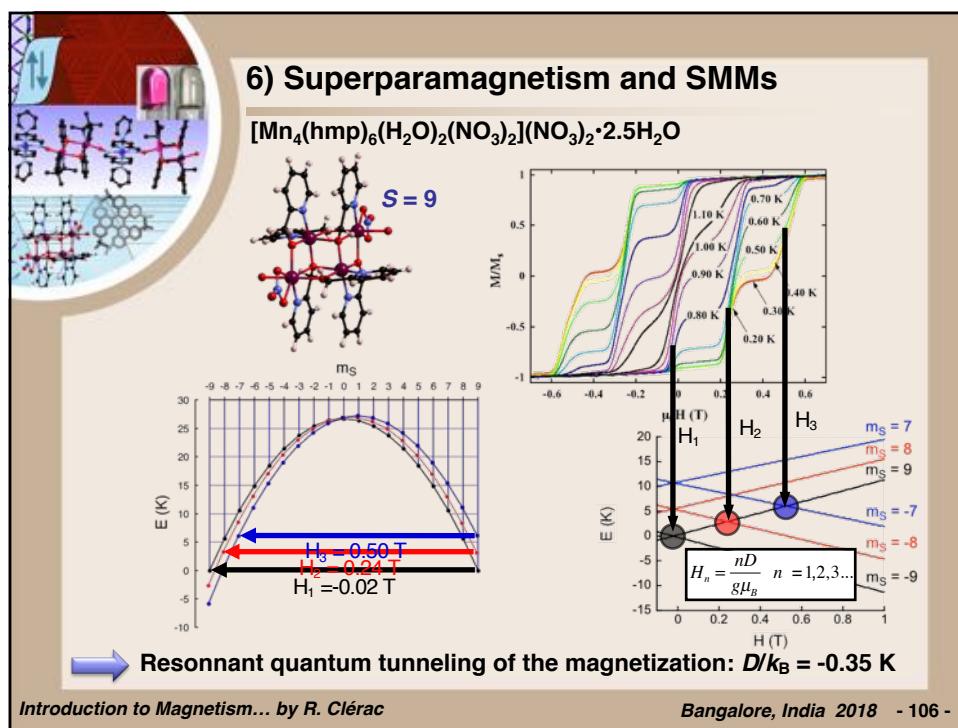
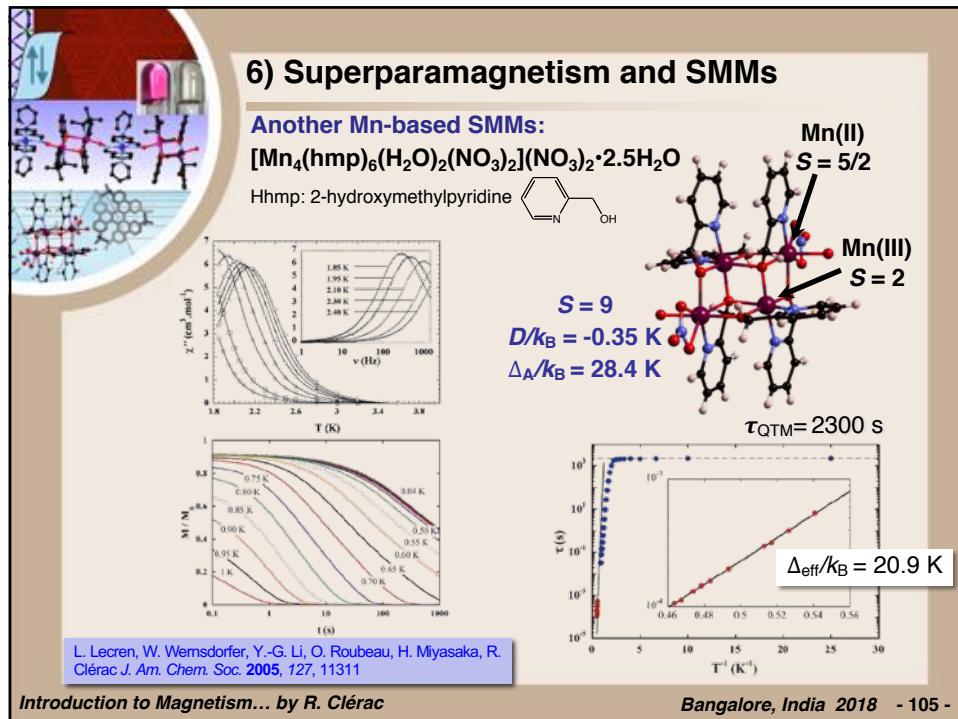
$$\Delta_{\text{eff}}/k_B \approx 61 \text{ K}$$



(a) Boyd, P. D. W.; Li, Q.; Vincent, J. B.; Folting, K.; Chang, H.-R.; Streib, W. E.; Huffman, J. C.; Christou, G.; Hendrickson, D. N. *J. Am. Chem. Soc.* **1988**, *110*, 8537; (b) Caneschi, A.; Gatteschi, D.; Sessoli, R. *J. Am. Chem. Soc.* **1991**, *113*, 5873; (c) Sessoli, R.; Tsai, H.-L.; Schake, A. R.; Wang, S.; Vincent, J. B.; Folting, K.; Gatteschi, D.; Christou, G.; Hendrickson, D. N. *J. Am. Chem. Soc.* **1993**, *115*, 1804; (d) Gatteschi, D.; Caneschi, A.; Pardi, L.; Sessoli, R. *Science* **1994**, *265*, 1054; Thomas, L.; Lioni, F.; Ballou, R.; Gatteschi, D.; Sessoli, R.; Barbara, B. *Nature* **1996**, *383*, 145.

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## Introduction to Magnetism

- 1) Introduction and definitions**
- 2) The Diamagnetism**
- 3) The Paramagnetism**
- 4) Magnetic phase transitions and magnetic orders**
- 5) Molecular magnetism**
- 6) Superparamagnetism and Single-Molecule Magnets**
- 7) Single-Chain Magnets**

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## 7) Single-Chain Magnets

**A real system made by design: a SMM building-block  
[Mn<sup>III</sup>(salen)-Fe<sup>III</sup>(CN)<sub>6</sub>-Mn<sup>III</sup>(salen)] trinuclear complexes:**

K[Mn<sub>2</sub>(5-Brsalen)<sub>2</sub>(H<sub>2</sub>O)<sub>2</sub>Fe(CN)<sub>6</sub>]·2H<sub>2</sub>O [1,4]

K[Mn<sub>2</sub>(5-Clsalen)<sub>2</sub>(H<sub>2</sub>O)<sub>2</sub>Fe(CN)<sub>6</sub>]·2H<sub>2</sub>O [1]

(NEt<sub>4</sub>)<sub>2</sub>[Mn<sub>2</sub>(5-Clsalen)<sub>2</sub>(H<sub>2</sub>O)<sub>2</sub>Fe(CN)<sub>6</sub>]·H<sub>2</sub>O [2]

(NEt<sub>4</sub>)<sub>2</sub>[Mn<sub>2</sub>(salmen)<sub>2</sub>(MeOH)<sub>2</sub>Fe(CN)<sub>6</sub>] [3,5]

$\text{X} = \text{H: salen}$   
 $\text{X} = \text{Br: 5-Brsalen}$   
 $\text{X} = \text{Cl: 5-Clsalen}$

$S = 9/2$   
 $D/k_B = -1.25 \text{ K}$

[1] H. Miyasaka, N. Matsumoto, H. Okawa, N. Re, E. Gallo, C. Floriani *J. Am. Chem. Soc.* **1996**, *118*, 981  
 [2] H. Miyasaka, N. Matsumoto, N. Re, E. Gallo, C. Floriani *Inorg. Chem.* **1997**, *36*, 670  
 [3] H. Miyasaka, H. Ieda, N. Matsumoto, N. Re, E. Crescenzi, C. Floriani *Inorg. Chem.* **1998**, *37*, 255  
 [4] H. J. Choi, J. J. Sokol, J. R. Long *Inorg. Chem.* **2004**, *43*, 1606  
 [5] M. Ferbinteanu, H. Miyasaka, W. Wernsdorfer, K. Nakata, K. Sugiura, M. Yamashita, C. Coulon, R. Clérac, *J. Am. Chem. Soc.* **2005**, *127*, 3090

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Bangalore, India 2018 - 108 -

**7) Single-Chain Magnets**

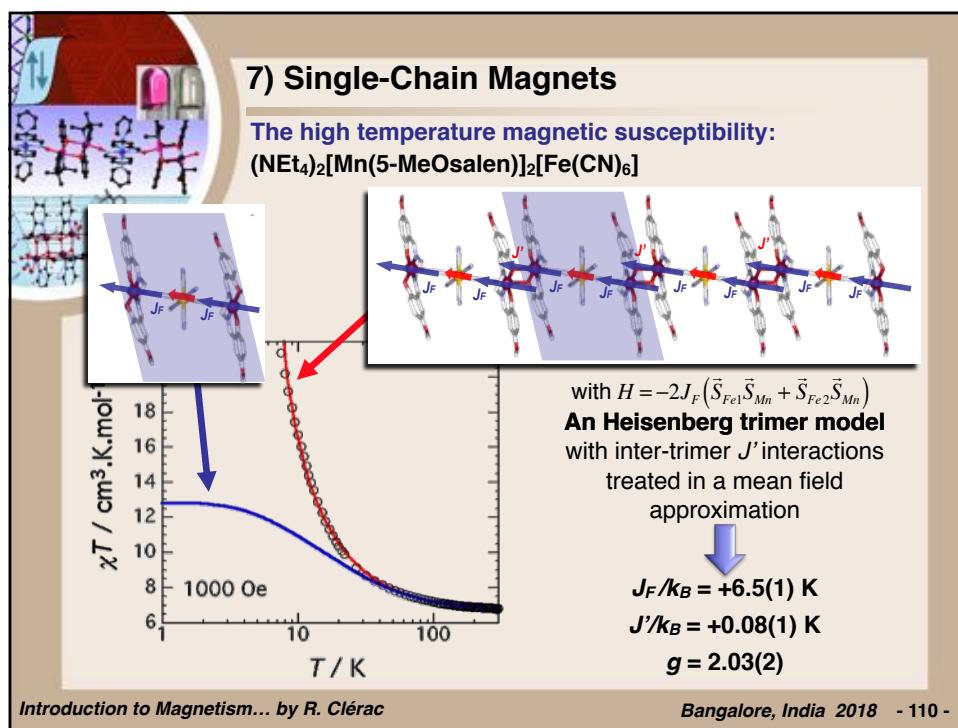
**The structural arrangement:**  
 $(\text{NEt}_4)_2[\text{Mn}(\text{5-MeOsalen})_2]\text{[Fe(CN)}_6]$

M. Ferbinteanu, H. Miyasaka, W. Wernsdorfer, K. Nakata, K. Sugiura, M. Yamashita, C. Coulon, R. Clérac, *J. Am. Chem. Soc.* 2005, 127, 3090

Repeating Unit

Isolated chains from a magnetic point of view

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**7) Single-Chain Magnets**

**A physicist view:**  
 $(\text{NEt}_4)_2[\text{Mn}(5\text{-MeOsalen})]_2[\text{Fe}(\text{CN})_6]$

$J_F/k_B = +6.5(1) \text{ K}$   
 $J'/k_B = +0.07(1) \text{ K}$   
 $J_{\text{Mn-Mn}}/k_B = +0.40(6) \text{ K}$

because  $|J_F| \gg J_{\text{Mn-Mn}}$  and  
 for  $|J_F| \gg k_B T$

$\left\langle S_T = 9/2 \right\rangle$   
 $S = 2, S = 1/2, S = 2$

because  $|J_F| \gg J_{\text{Mn-Mn}}$  and  
 for  $|J_F| \gg k_B T$

Chain of ferromagnetically coupled anisotropic  $S = 9/2$  spins

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**7) Single-Chain Magnets**

**Single crystal measurements:  $H$  in the hard plane**  
 $(\text{NEt}_4)_2[\text{Mn}(5\text{-MeOsalen})]_2[\text{Fe}(\text{CN})_6]$

$M = -2J' \sum_{-\infty}^{+\infty} \vec{S}_{T,i} \vec{S}_{T,i+1} + D \sum_{-\infty}^{+\infty} \vec{S}_{T,i}^2$

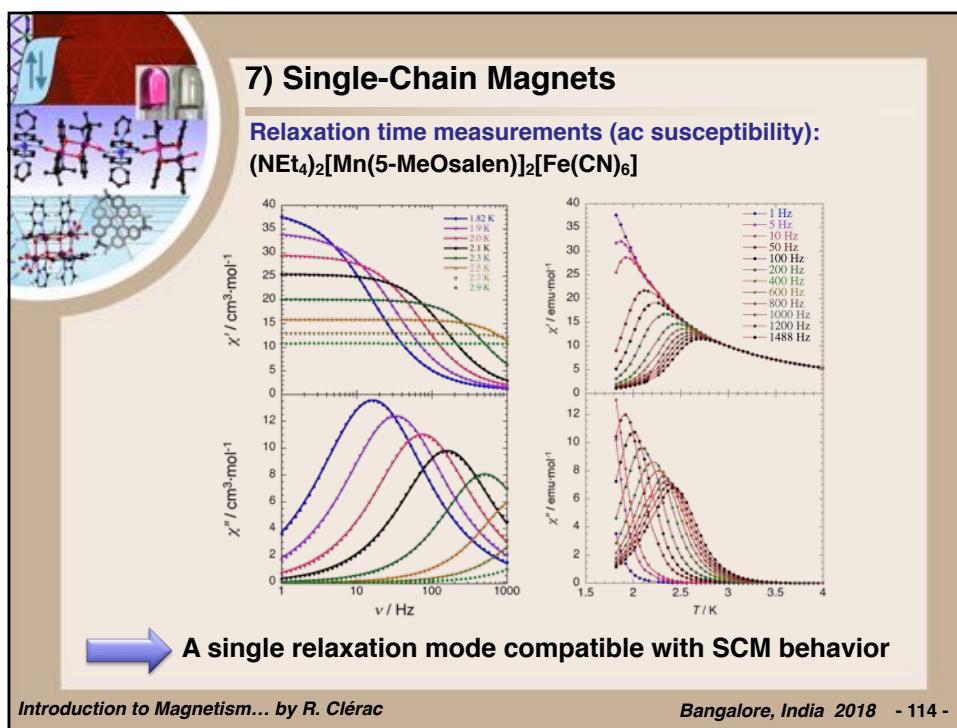
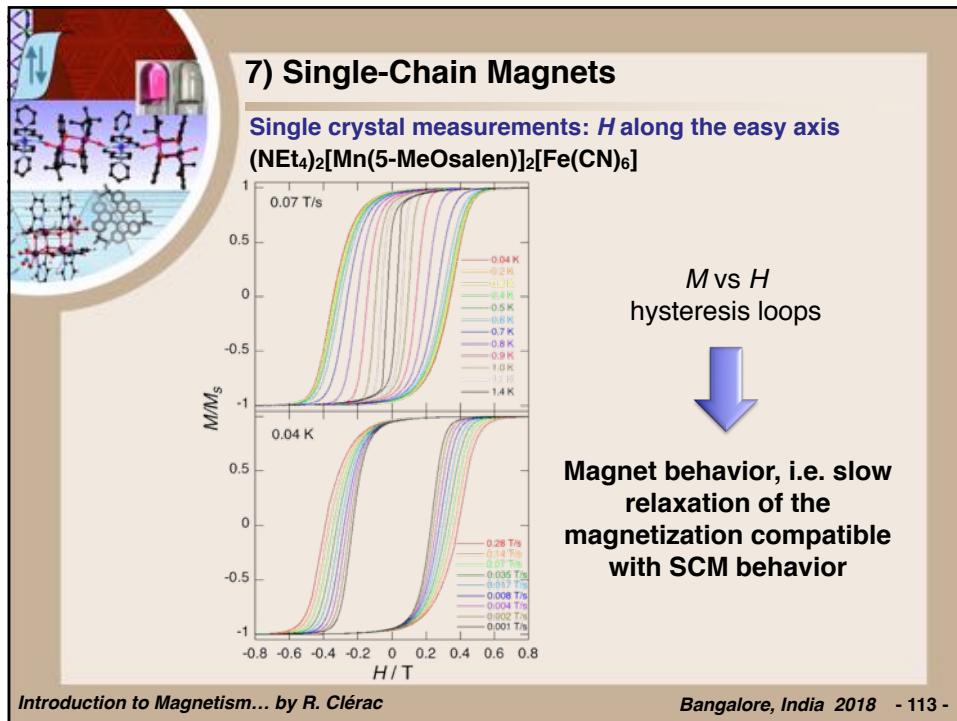
Estimation of  $D$

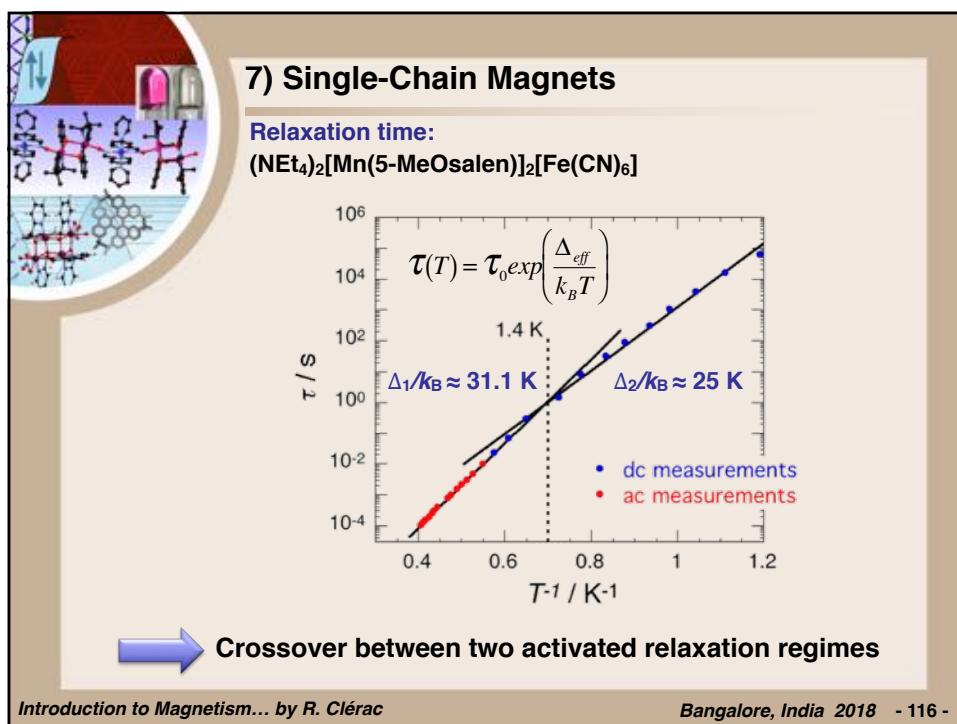
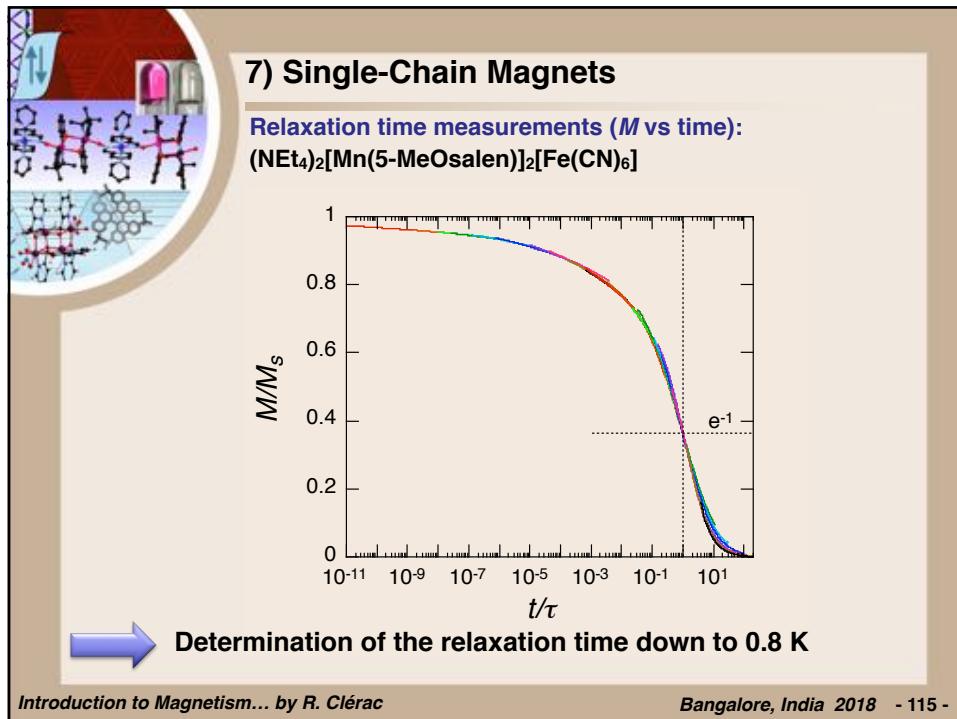
$2DS_T^2 \approx g\mu_B S_T H_a$

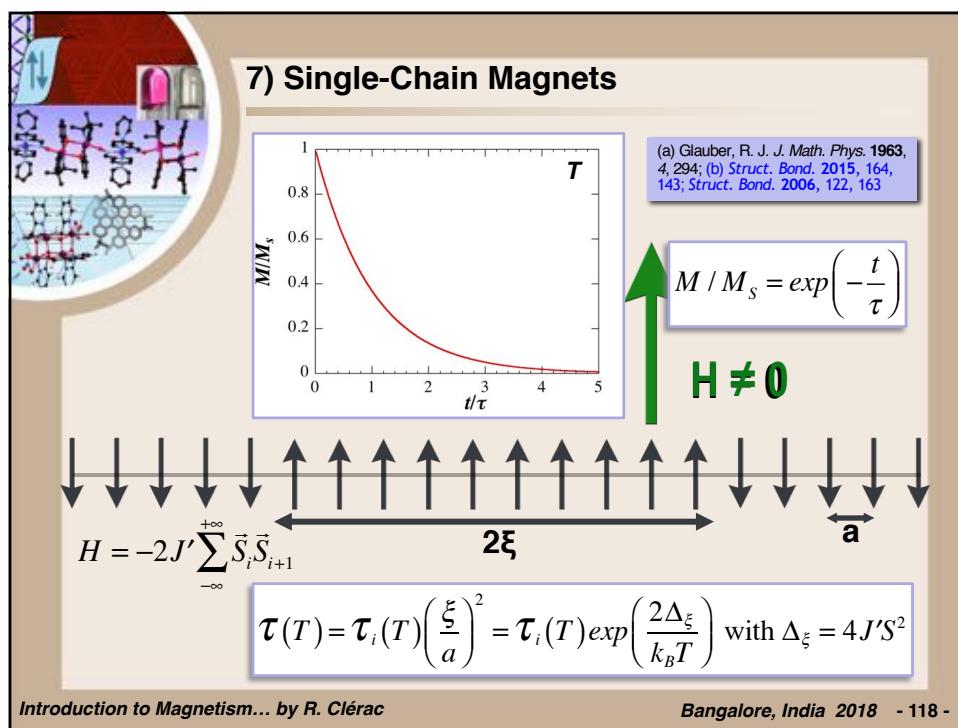
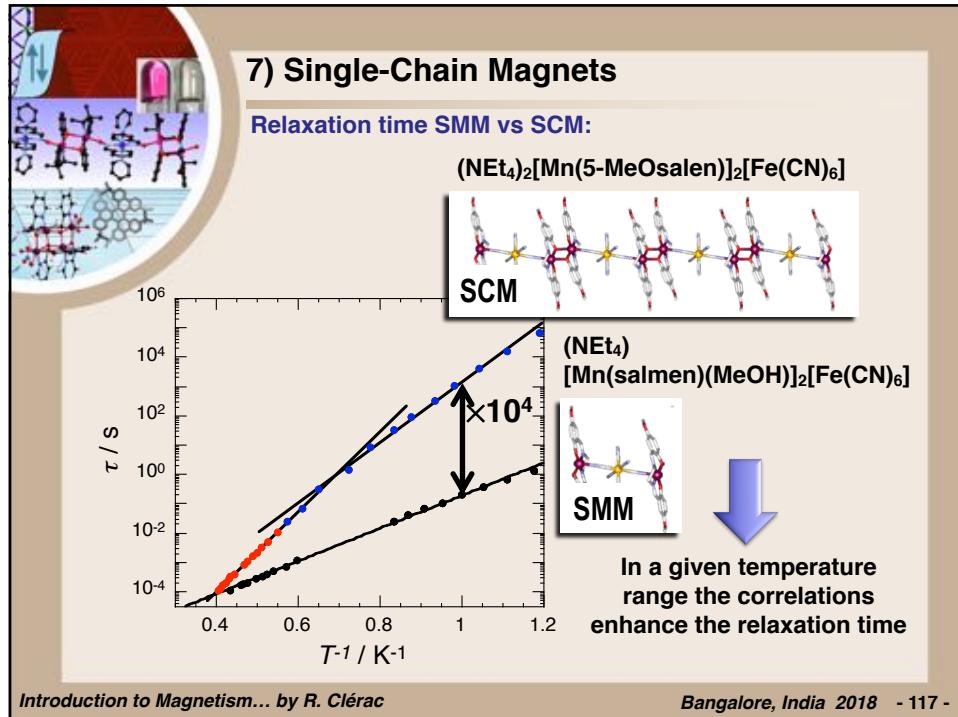
$\downarrow$   
 $D/k_B \approx -0.94 \text{ K}$

$D/k_B \approx -1.25 \text{ K}$

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**7) Single-Chain Magnets**

**The relaxation time in a Single-Chain Magnet:  
Chain of ferromagnetically coupled Ising spins (Glauber)**

$$\tau(T) = \tau_i(T) \exp\left(\frac{2\Delta_\xi}{k_B T}\right) \text{ with } \Delta_\xi = 4J'S^2$$

(a) Glauber, R. J. *J. Math. Phys.* 1963, 4, 294; (b) *Struct. Bond.* 2015, 164, 143; *Struct. Bond.* 2006, 122, 163

1) For a chain of anisotropic spins ( $D < 0$  and  $|D/J| > 4/3$ : Ising limit) :

$$H = -2J' \sum_{-\infty}^{+\infty} \vec{S}_{T,i} \vec{S}_{T,i+1} + D \sum_{-\infty}^{+\infty} \vec{S}_{T,iz}^2$$

$$\tau_i(T) = \tau_0 \exp\left(\frac{\Delta_A}{k_B T}\right) \text{ with } \Delta_A = |D|S_T^2 \quad \rightarrow \quad \tau(T) = \tau_0 \exp\left(\frac{2\Delta_\xi + \Delta_A}{k_B T}\right)$$

Coulon, C.; Clérac, R.; Lecren, L.; Wernsdorfer, W.; Miyasaka, H. *Phys. Rev. B* 2004, 69, 132408

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**7) Single-Chain Magnets**

**Relaxation time: the Glauber's regime  
(NEt<sub>4</sub>)<sub>2</sub>[Mn(5-MeOsalen)]<sub>2</sub>[Fe(CN)<sub>6</sub>]**

with  $S_T = 9/2$   
 $J'/k_B \approx 0.08 \text{ K}$   
 $D/k_B = -0.94 \text{ K}$

$$\tau(T) = \tau_0 \exp\left(\frac{2\Delta_\xi + \Delta_A}{k_B T}\right)$$

with

$$2\Delta_\xi + \Delta_A = 8J'S_T^2 + |D|S_T^2$$

$\Delta_1/k_B \approx 31.1 \text{ K}$   
 $\Delta_2/k_B \approx 25 \text{ K}$

dc measurements (blue dots)  
ac measurements (red dots)

$\Delta_1/k_B = (8J' + |D|)S_T^2/k_B = 32 \text{ K}$

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**7) Single-Chain Magnets**

**The relaxation time in a Single-Chain Magnet:**

**2) For a chain of anisotropic spins with a length  $L$  ( $D < 0$  and  $|D/J| > 4/3$  Ising limit):**

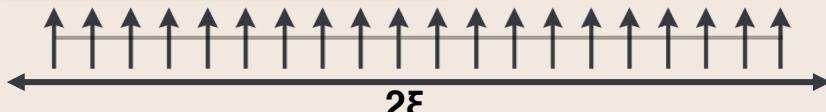
$$H = -2J \sum_{-\infty}^{+\infty} \vec{S}_{T,i} \vec{S}_{T,i+1} + D \sum_{-\infty}^{+\infty} \vec{S}_{T,i}^2$$

(a) Coulon, C.; Miyasaka, H.; Clérac, R. *Struct. Bond.* 2006, 122, 163  
(b) Coulon, C.; Clérac, R.; Lecren, L.; Wernsdorfer, W.; Miyasaka, H. *Phys. Rev. B* 2004, 69, 132408

for  $2\xi < L$   $\rightarrow \tau(T) = \tau_0 \exp\left(\frac{2\Delta_\xi + \Delta_A}{k_B T}\right)$

for  $2\xi > L$   $\rightarrow \tau(T) = \tau_0 \exp\left(\frac{\Delta_\xi + \Delta_A}{k_B T}\right)$

Luscombe, J. H. et al., *Phys. Rev. E* 1996, 53, 5852  
Leal da Silva, J. K. et al., *Phys. Rev. E* 1995, 52, 4527



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**7) Single-Chain Magnets**

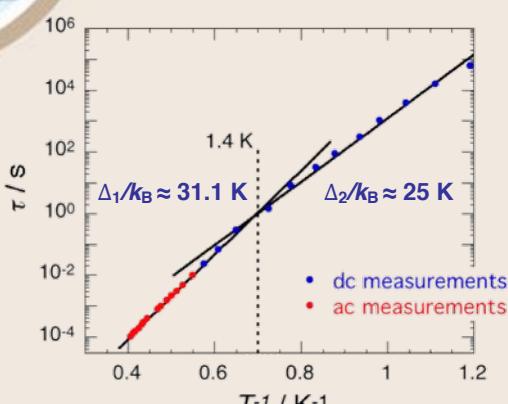
**Relaxation time: the finite-size chain regime**  
 $(\text{NEt}_4)_2[\text{Mn}(5\text{-MeOsalen})]_2[\text{Fe}(\text{CN})_6]$

with  $S_T = 9/2$   
 $J/k_B \approx 0.08 \text{ K}$   
 $D/k_B = -0.94 \text{ K}$

$\tau(T) = \tau_0 \exp\left(\frac{\Delta_\xi + \Delta_A}{k_B T}\right)$   
with  
 $\Delta_\xi + \Delta_A = 4J'S_T^2 + |D|S_T^2$

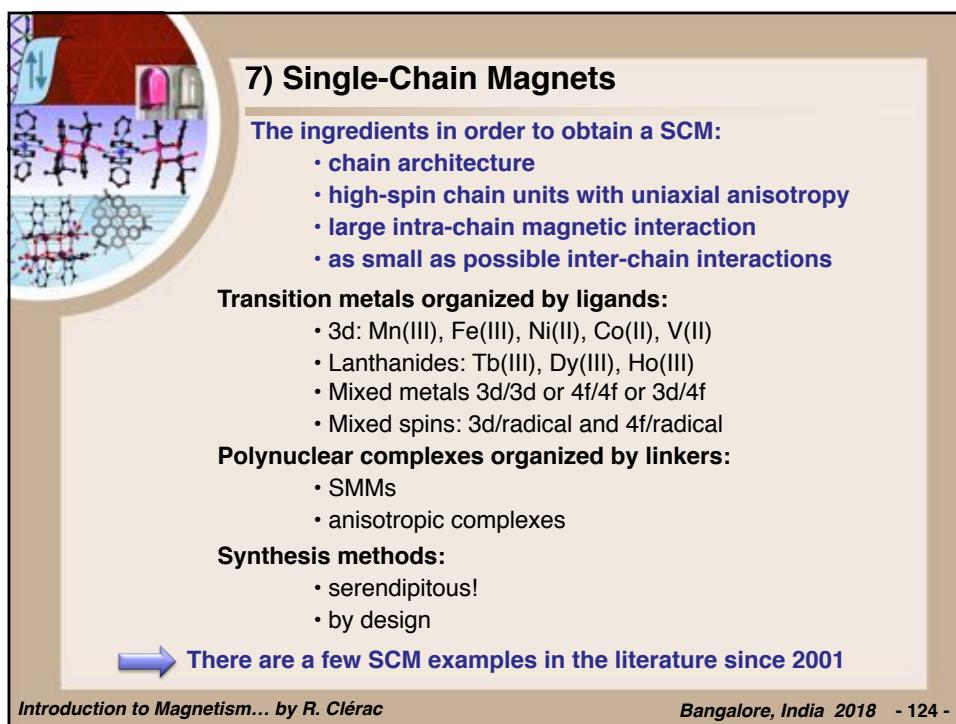
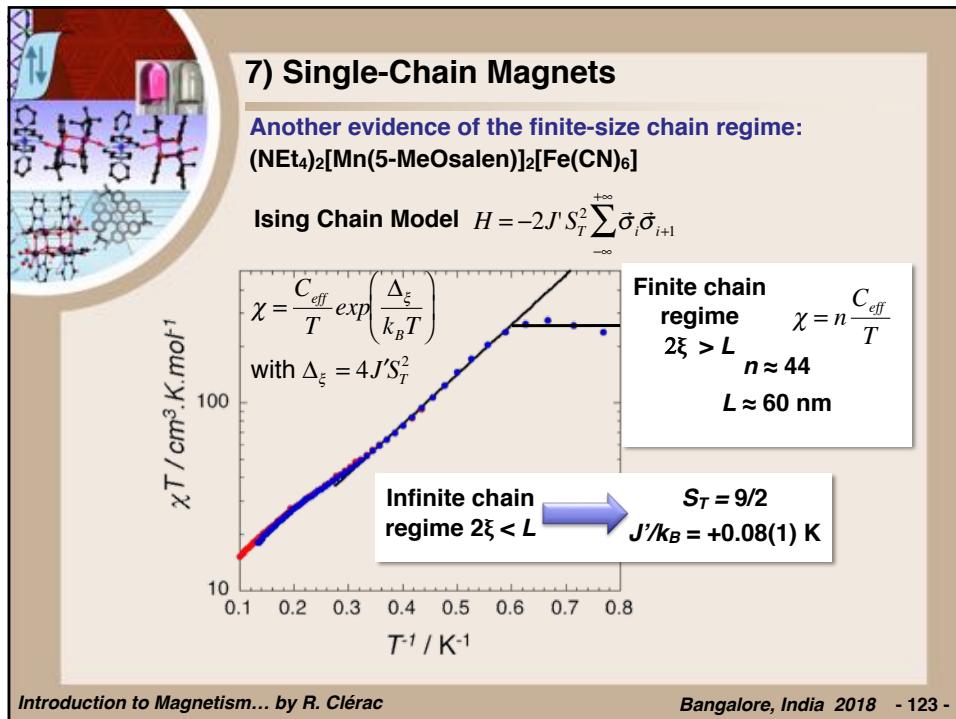
$\Delta_1/k_B \approx 31.1 \text{ K}$   
 $\Delta_2/k_B \approx 25 \text{ K}$

$\bullet$  dc measurements  
 $\bullet$  ac measurements



$\Delta_2/k_B = (4J' + |D|)S_T^2/k_B = 25.5 \text{ K}$

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**7) Single-Chain Magnets**

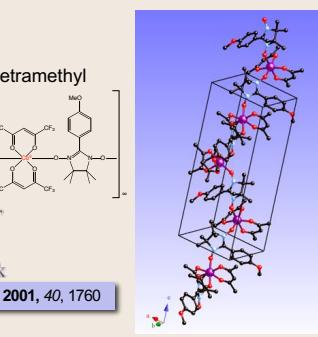
**The “first” chain exhibiting slow relaxation:**  
**Co<sup>II</sup>(hfac)<sub>2</sub>(NITPhOMe)**

hfac: hexafluoroacetylacetone  
 NITPhOMe: 4'-methoxy-phenyl-4,4,5,5-tetramethyl imidazoline-1-oxyl-3-oxide

**Cobalt(II)-Nitronyl Nitroxide Chains as Molecular Magnetic Nanowires**

Andrea Caneschi, Dante Gatteschi,\* Nikolia Lalioti, Claudio Sangregorio, Roberta Sessoli, Giovanni Venturi, Alessandro Vindigni, Angelo Rettori, Maria G. Pini, and Miguel A. Novak

*Angew. Chem. Int. Ed.* 2001, 40, 1760



Three plots showing magnetic properties: 1) Hysteresis loop (M vs H) with  $H_c \approx 1$  T; 2) Temperature dependence of magnetization ( $\chi_M$ ) with peaks at low temperatures; 3) Arrhenius plot of ln(t/Sec) vs 1/T (K<sup>-1</sup>) showing a linear relationship with  $\Delta_{eff}/k_B = 153$  K.

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**7) Single-Chain Magnets**

**The first chain named “Single-Chain Magnet”:**  
**[Mn<sub>2</sub>(saltmen)<sub>2</sub>Ni(pao)<sub>2</sub>(py)<sub>2</sub>](ClO<sub>4</sub>)<sub>2</sub>**

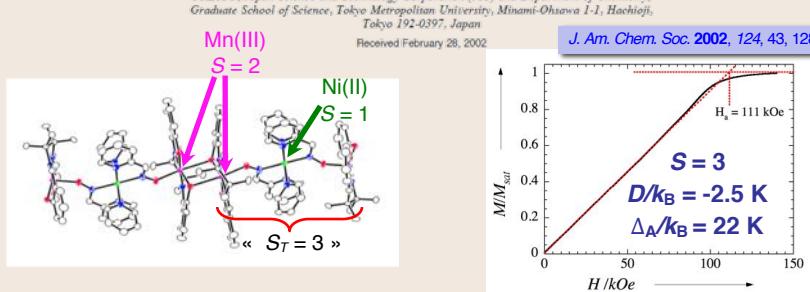
saltmen<sup>2-</sup>: N,N'-(1,1,2,2-tetramethylethylen) bis(salicylideneiminate)  
 pao: pyridine-2-aldoximate; py: pyridine

**Evidence for Single-Chain Magnet Behavior in a Mn<sup>III</sup>–Ni<sup>II</sup> Chain Designed with High Spin Magnetic Units: A Route to High Temperature Metastable Magnets**

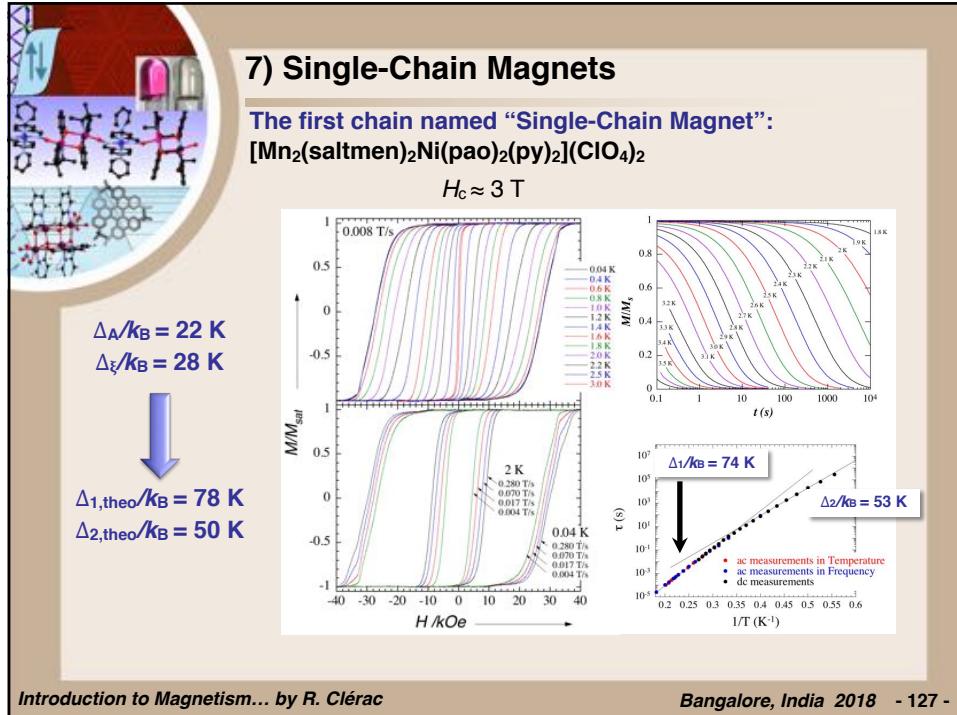
Rodolphe Clérac,<sup>5,†</sup> Hitoshi Miyasaka,<sup>4,5,§</sup> Masahiro Yamashita,<sup>4</sup> and Claude Coulon<sup>7</sup>

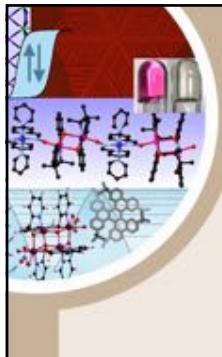
Contribution from the Centre de Recherche Paul Pascal, CNRS UPR 8641, avenue du Dr. A. Schweitzer, 33600 Pessac, France, and “Structural Ordering and Physical Properties”, PRESTO, Japan Science and Technology Corporation (JST) and Department of Chemistry, Graduate School of Science, Tokyo Metropolitan University, Minami-Ohsawa 1-1, Hachioji, Tokyo 192-0397, Japan

Received February 28, 2002      *J. Am. Chem. Soc.* 2002, 124, 43, 12837



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The end...

# Thank you for Your attention!

*clerac@crpp-bordeaux.cnrs.fr*

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*Bangalore, India 2018 - 129 -*